

Spectral universality for a general class of matrices

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37th Conference on Stochastic Processes and Applications
Buenos Aires, Aug 1, 2014

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Partially supported by ERC Advanced Grant, RANMAT, No. 338804

INTRODUCTION

Basic question [Wigner]: What can be said about the statistical properties of the eigenvalues of a large random matrix? Do some universal patterns emerge?

$$H = \begin{pmatrix} h_{11} & h_{12} & \dots & h_{1N} \\ h_{21} & h_{22} & \dots & h_{2N} \\ \vdots & \vdots & & \vdots \\ h_{N1} & h_{N2} & \dots & h_{NN} \end{pmatrix} \implies (\lambda_1, \lambda_2, \dots, \lambda_N) \text{ eigenvalues?}$$

N = size of the matrix, will go to infinity.

Analogy: Central limit theorem: $\frac{1}{\sqrt{N}}(X_1 + X_2 + \dots + X_N) \sim \mathcal{N}(0, \sigma^2)$

Wigner Ensemble:

$H = (h_{jk})_{1 \leq j, k \leq N}$ **complex hermitian** $N \times N$ matrix

$h_{jk} = \bar{h}_{kj}$ (for $j < k$) are complex and h_{kk} are real independent random variables with normalization

$$\mathbb{E}h_{jk} = 0, \quad \mathbb{E}|h_{jk}|^2 = \frac{1}{N}.$$

The eigenvalues $\lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_N$ are **of order one**: (on average)

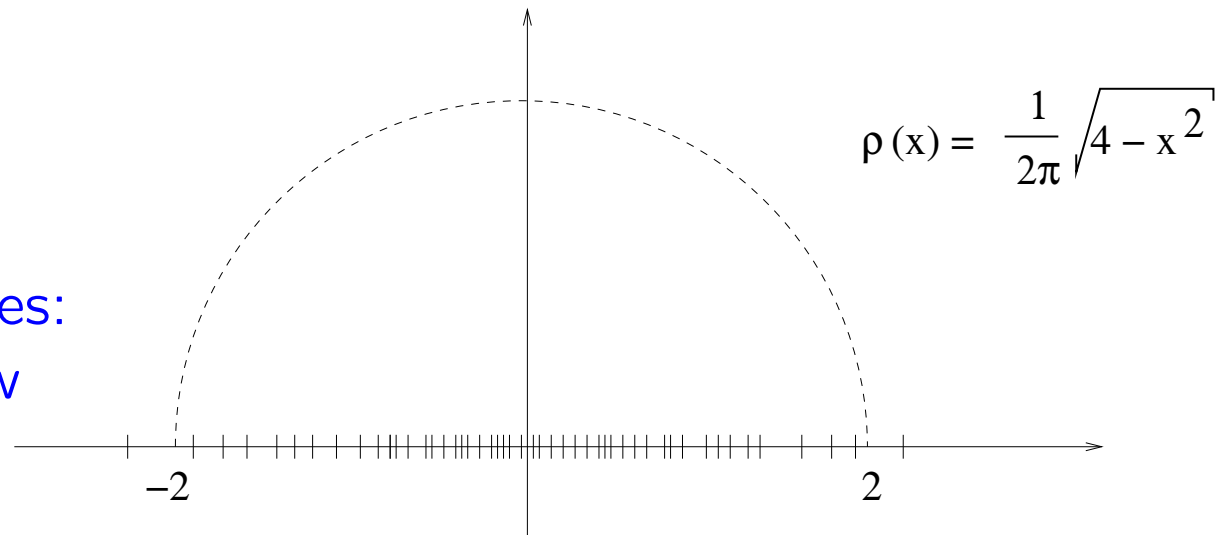
$$\mathbb{E} \frac{1}{N} \sum_i \lambda_i^2 = \mathbb{E} \frac{1}{N} \text{Tr} H^2 = \frac{1}{N} \sum_{ij} \mathbb{E}|h_{ij}|^2 = 1$$

Complex hermitian can be replaced with real symmetric or quaternion self-dual.

If h_{ij} is Gaussian, then GUE, GOE, GSE.

Wigner's observations (holds for all symmetry classes)

i) Density of eigenvalues:
Wigner semicircle law



ii) Level repulsion: Wigner surmise (in the bulk and for GOE)

$$\mathbb{P}\left(N(\lambda_{i+1} - \lambda_i) = s + ds\right) \approx \frac{\pi s}{2} \exp\left(-\frac{\pi}{4}s^2\right) ds$$

Guessed by a 2x2 matrix calculation

SINE KERNEL FOR CORRELATION FUNCTIONS

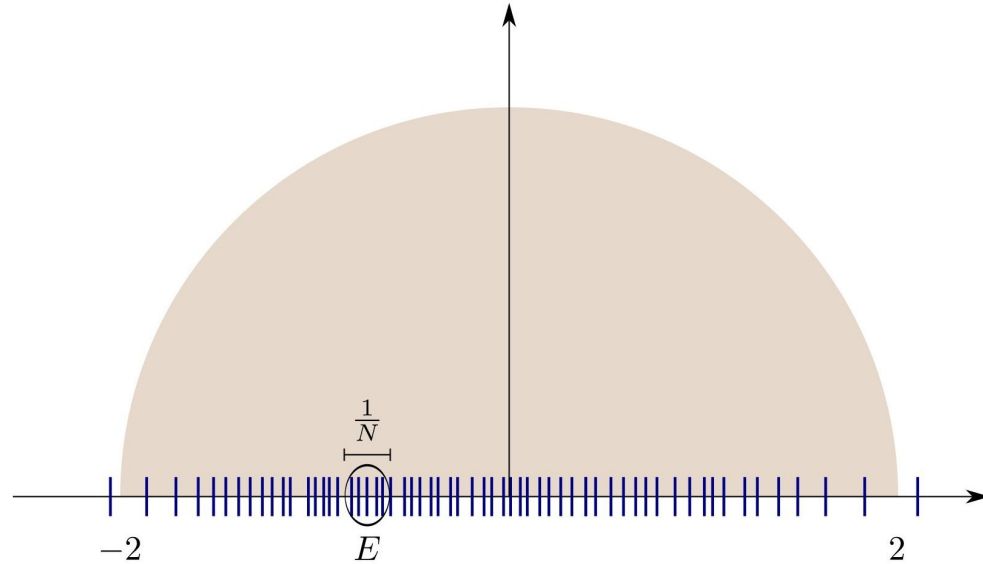
Probability density of the eigenvalues: $p(x_1, x_2, \dots, x_N)$

The k -point correlation function is given by

$$p_N^{(k)}(x_1, x_2, \dots, x_k) := \int_{\mathbb{R}^{N-k}} p(x_1, \dots, x_k, x_{k+1}, \dots, x_N) dx_{k+1} \dots dx_N$$

Special case: $k = 1$ (density)

$$\varrho_N(x) := p_N^{(1)}(x) = \int_{\mathbb{R}^{N-1}} p(x, x_2, \dots, x_N) dx_2 \dots dx_N$$



Rescaled correlation functions at energy E

$$p_E^{(k)}(\mathbf{x}) := \frac{1}{[\varrho(E)]^k} p_N^{(k)}\left(E + \frac{x_1}{N\varrho(E)}, E + \frac{x_2}{N\varrho(E)}, \dots, E + \frac{x_k}{N\varrho(E)}\right)$$

Rescales the gap $\lambda_{i+1} - \lambda_i$ to $O(1)$.

Local level correlation statistics for GUE [Gaudin, Dyson, Mehta]

k -point correlation functions are given by $k \times k$ determinants:

$$\lim_{N \rightarrow \infty} p_E^{(k)}(\mathbf{x}) = \det \left\{ S(x_i - x_j) \right\}_{i,j=1}^k, \quad S(x) := \frac{\sin \pi x}{\pi x}$$

The limit is independent of E as long as $|E| < 2$ (bulk spectrum)

Gap distribution can be obtained from correlation functions by the exclusion-inclusion formula. Wigner surmise is quite precise.

Main question: going beyond Gaussian towards universality!

Wigner-Dyson-Mehta conjecture: Local statistics is universal in the bulk spectrum for any Wigner matrix; only symmetry type matters.

Solved recently for any symmetry class:

[E-Schlein-Peche-Ramirez-Yau, 2009] – Hermitian case, fixed E

[E-Schlein-Yau-Yin, 2010] – averaged E

[E-Yau, 2012] – fixed gap label

[Bourgade-E-Yau-Yin, 2014] – fixed E

Related results:

[Johansson, 2000] Hermitian case with large Gaussian components

[Tao-Vu, 2009] Hermitian case via moment matching.

(Similar development for the edge and for β -log gases).

Three-step strategy:

1. Local (entry-wise) semicircle law down to scales $\gg 1/N$.
2. Use local equilibration of Dyson Brownian motion to prove universality for matrices with a tiny Gaussian component
3. Use perturbation theory to remove the tiny Gaussian component.

All these results were obtained under the condition that the matrix elements h_{ij} are independent, centered and

$$\sum_j s_{ij} = 1, \quad s_{ij} := \mathbb{E}|h_{ij}|^2$$

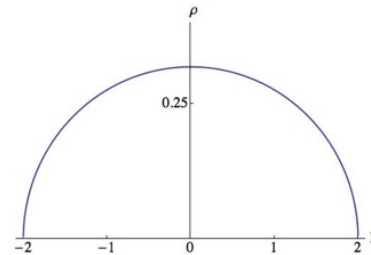
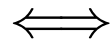
Today's talk is about Wigner matrices **without this red** condition. We'll call them **Wigner-type** matrices.

Red was used at many places in the previous analysis:

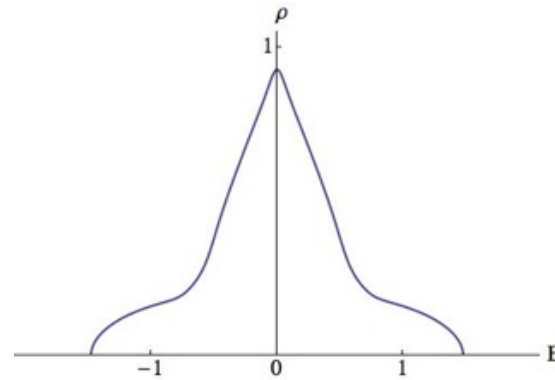
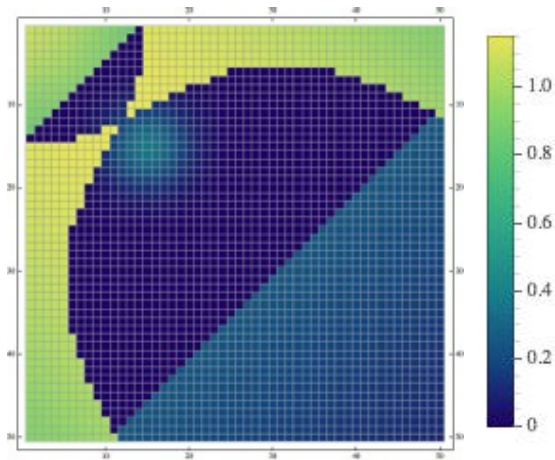
- Limiting density is explicit.
- Homogeneity: $G_{ii} \approx G_{jj}$ for the resolvent matrix elements.
- DBM: Initial data is already close to global equilibrium.

Variance profile and limiting density of states (DOS)

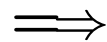
$$\sum_j s_{ij} = 1$$



General variance profile $s_{ij} = \mathbb{E}|h_{ij}|^2$: not the semicircle any more.



$$\sum_j s_{ij} \neq \text{const}$$



Density of states

Main theorem (informally)

Theorem [Ajanki-E-Krüger] Let $H = H^*$ be a Wigner-type matrix

$$\bar{h}_{ji} = h_{ij} \quad \text{independent, centered}$$

$$\mathbb{E}|h_{ij}|^2 = s_{ij} = \frac{1}{N} S\left(\frac{i}{N}, \frac{j}{N}\right)$$

with a limiting profile function $S : [0, 1]^2 \rightarrow \mathbb{R}_+$. Then for the matrix elements of the resolvent $G = (H - z)^{-1}$, we have

$$G_{ij}(z) \approx \delta_{ij} m_{\frac{i}{N}}(z)$$

where $m_x(z)$ solves the self-consistent equation

$$-\frac{1}{m_x(z)} = z + \int_0^1 S(x, y) m_y(z) dy \quad (*)$$

Limiting DOS $\rho(E) := \frac{1}{\pi} \int_0^1 \text{Im } m_x(E + i0) dx$

Note: The nonlinear **vector** equation (*) replaces the usual self consistent **scalar** eq $m^{-1} = -(z + m)$ of the semicircle density.

Constantness of row sums, $\sum_j s_{ij} = 1$, implies semicircle

$$G(z) := (H - z)^{-1}, \quad z = E + i\eta, \quad \eta > 0$$

Schur formula with the resolvent of the i -th minor

$$\begin{aligned} \frac{1}{G_{ii}} &= h_{ii} - z - \sum_{ab} h_{ia} G_{ab}^{(i)} h_{bi} \\ &\approx -z - \mathbb{E}^{(i)} \sum_{ab} h_{ia} G_{ab}^{(i)} h_{bi} \\ &\approx -z - \sum_a s_{ia} G_{aa} \end{aligned}$$

Fact: The self-consistent equation

$$\frac{1}{m_i} = -z - \sum_a s_{ia} m_a, \quad \text{Im } m_a > 0,$$

has a unique solution.

It is constant, $m_a = m$, **iff** the row sums are constant and then

$$\frac{1}{m} = -z - m \quad \implies \quad \text{Stieltjes transform of the semicircle}$$

Quadratic vector equation (QVE)

Suppose s_{ij} is given by a limiting profile function $S : [0, 1]^2 \rightarrow \mathbb{R}_+$:

$$s_{ij} = \frac{1}{N} S\left(\frac{i}{N}, \frac{j}{N}\right),$$

Continuum limit of the self-consistent equation for $G_{ii}(z) \approx m_{\frac{i}{N}}(z)$

$$-\frac{1}{m(z)} = z + Sm(z), \quad (Sf)_x = \int_0^1 S(x, y) f_y dy \quad (\text{QVE})$$

For any $z \in \mathbb{H}$ (complex upper half plane), we consider solutions under the constraint $\text{Im } m > 0$,

Fact: Solution exists and is unique.

Fact: The solution is not constant in general. Semicircle is the “easiest” case.

Two main steps for the proof of the main theorem:

- 1) Analyse the solution of the continuum QVE, including its stability (no matrix, no N)
- 2) Prove: the resolvent of the RM is close to the solution of QVE. (Schur, fluctuation averaging, dichotomy becomes trichotomy)

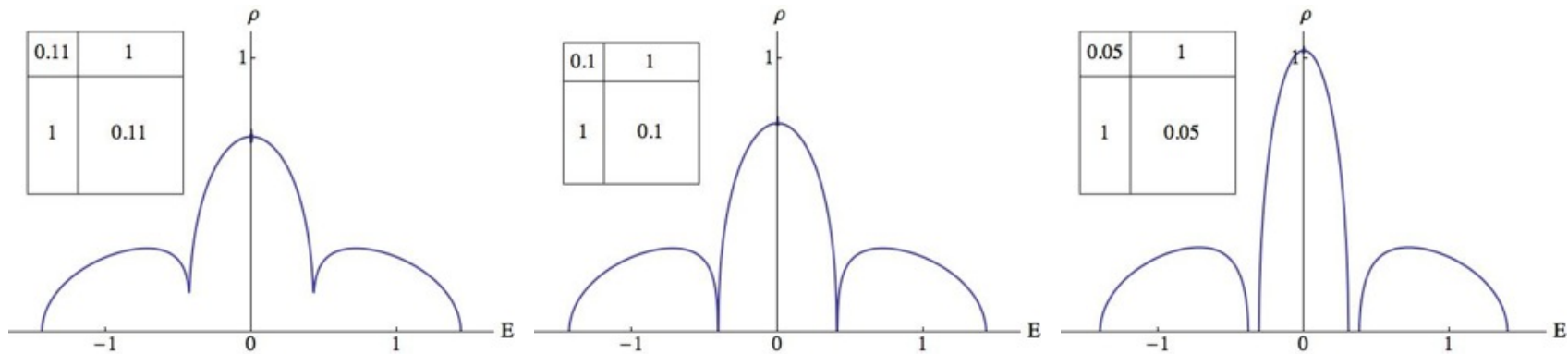
Note: If $\sum_j s_{ij} = 1$, Step 1 is trivial, since the solution $m_x(z) = m(z)$ is given explicitly by a quadratic **scalar** equation $-m^{-1} = z + m$. So all previous efforts to prove local semicircle law was in Step 2.

If $\sum_j s_{ij} \neq \text{const}$, Step 1 is nontrivial and gives rise to a complex pattern.

Despite its natural form, QVE has not been studied quantitatively.

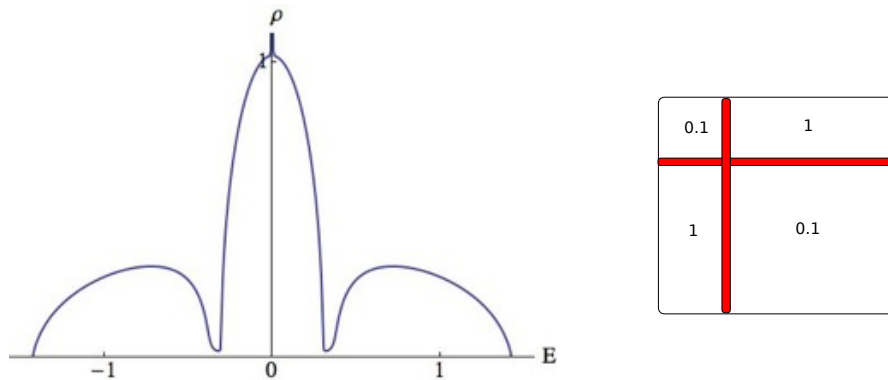
Features of the DOS for Wigner-type matrices

1) Support splits via cusps:



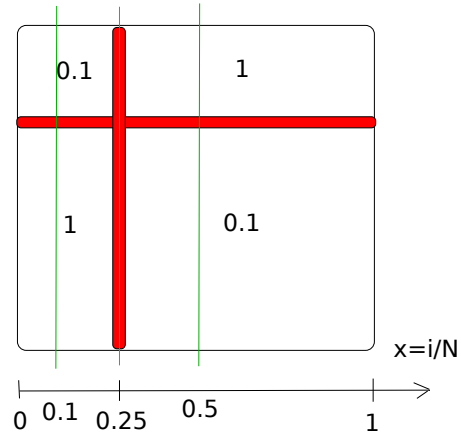
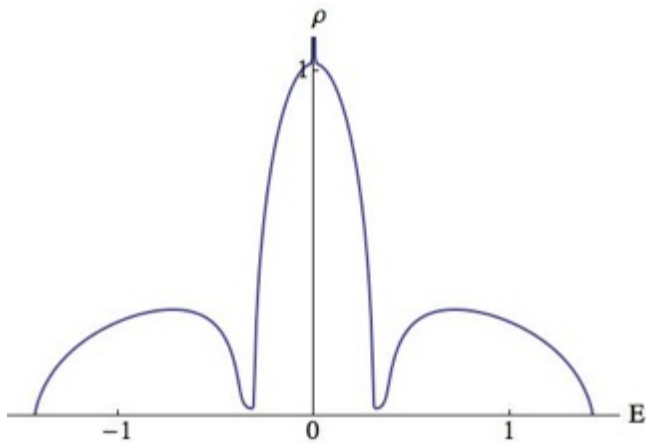
(Matrices in the pictures represent the variance matrix)

2) Smoothing of the S -profile avoids splitting (\Rightarrow single interval)



DOS of the same matrix as above but discontinuities in S are regularized

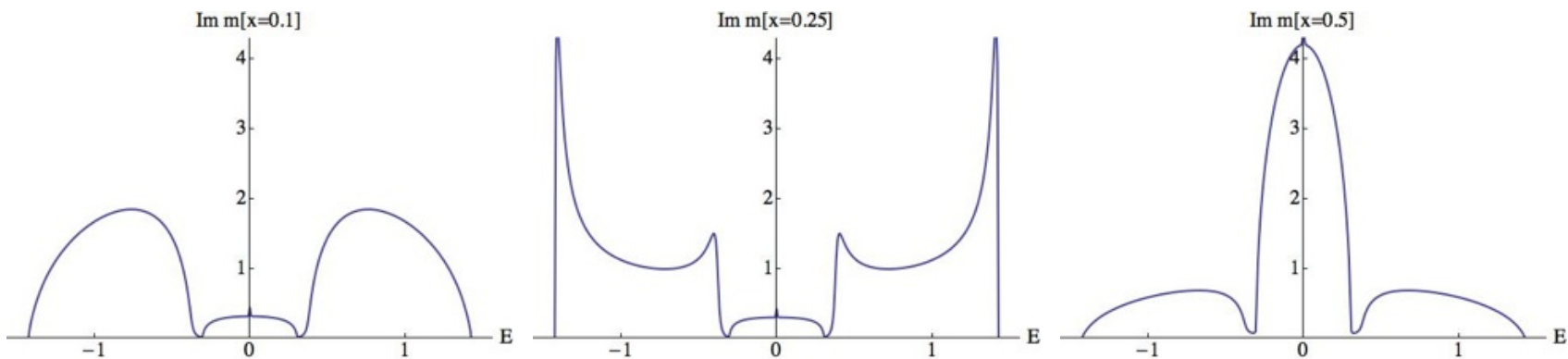
Relation between m_x and $m := \mathbf{A}v_x m_x$



$x = \frac{i}{N}$ continuum coordinates

Red: some interpolation

$\text{Im } m_x \not\approx \rho = \text{Im } m$. It may even behave very differently for some x :



Sections of $\text{Im } m_x(E)$ at various x 's indicated by the green lines.

Natural questions

- 1) How many intervals are there and what determines them?
- 2) Blow-up features and instability mechanisms
- 3) Universality of the singularity patterns?

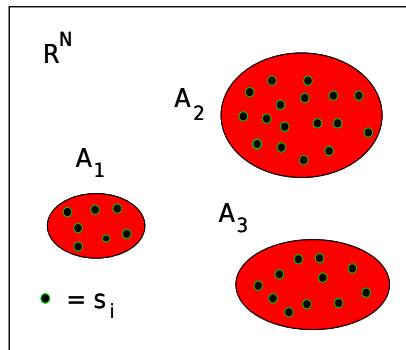
Number of intervals in the support of the DOS

Consider the set of row vectors of S

$$A := \{s_i : i = 1, 2, \dots, N\} \subset \mathbb{R}^N, \quad (s_i)_j := s_{ij}$$

Partition

$$A = A_1 \cup A_2 \cup \dots \cup A_n, \quad \text{s.t.} \quad \text{dist}(A_k, A_\ell) \geq \delta$$



Key object:
 n = Number of lumps

Conjecture: # spectral intervals $\leq 2n - 1$. We proved for $n = 1$

E.g. $s_{ij} = \frac{1}{N} S(\frac{i}{N}, \frac{j}{N})$ with $S(x, y)$ smooth $\Rightarrow n = 1$

Theorem [Ajanki-E-Krüger] If all lumps are macroscopic in the sense

$$\inf_x \int_0^1 \frac{1}{\|s_x - s_y\|^2} dy \geq C > 0$$

then the solution $m_x(z)$ of

$$-\frac{1}{m} = z + Sm, \quad (QVE)$$

is bounded; $m_x(z)$ is the Stieltjes transform of an a.c. measure

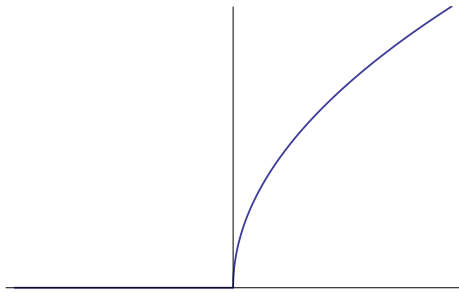
$$m_x(z) = \int_{\mathbb{R}} \frac{v_x(s)}{s - z} ds.$$

If S is irreducible, then the components are comparable, $\frac{v_x(E)}{v_y(E)} \sim 1$ for all E . The density, $\varrho = \int v_x dx$, is compactly supported, bounded, and it has a **universal** shape near the points when it (almost) vanishes.

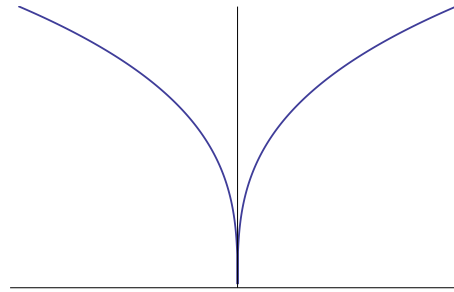
In particular, blowup can occur only if there is a small lump.

(Discontinuity in S is OK, but isolated row is not)

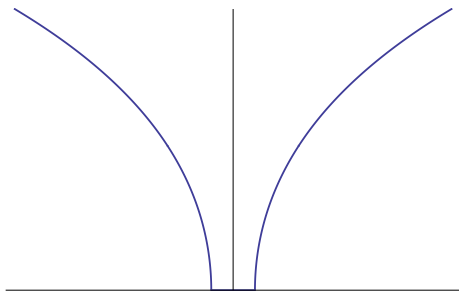
Universality of the singularities in the DOS



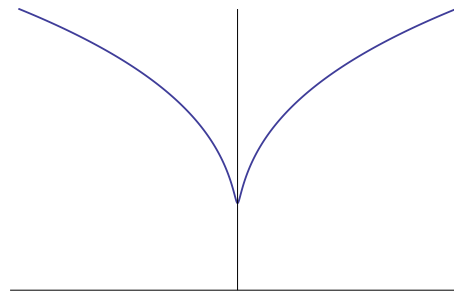
Edge, \sqrt{E} singularity



Cusp, $|E|^{1/3}$ singularity



Small-gap



Smoothed cusp

$$\frac{(2+\tau)\tau}{1+(1+\tau+\sqrt{(2+\tau)\tau})^{2/3}+(1+\tau-\sqrt{(2+\tau)\tau})^{2/3}}$$

$$\tau := \frac{|E|}{\text{gap}},$$

$$\frac{\sqrt{1+\tau^2}}{(\sqrt{1+\tau^2}+\tau)^{2/3}+(\sqrt{1+\tau^2}-\tau)^{2/3}-1} - 1$$

$$\tau := \frac{|E|}{(\text{minimum})^{1/3}}$$

Why cubic?

Stability of QVE: (used for both QVE analysis and RM)

$$-\frac{1}{m} = z + Sm, \quad -\frac{1}{m^{(\varepsilon)}} = z + Sm^{(\varepsilon)} + \varepsilon, \quad \|\varepsilon\| \ll 1$$

Decompose along the evector $(1 - |m|^2 S)f = 0$ ($f > 0$)

$$m^{(\varepsilon)} - m = \Theta f + v, \quad \|v\| \leq \Theta^2 + O(\varepsilon)$$

Decompose the third order perturbation expansion along f and f^\perp .

$$\tau_3 \Theta^3 + \tau_2 \Theta^2 + \tau_1 \Theta \sim O(\varepsilon)$$

Facts

$$|\tau_3| + |\tau_2| \neq 0 \quad \implies \quad \text{not more than cubic}$$

$$\tau_2 = \langle (\text{sgnRe } m), f^3 \rangle$$

If $\tau_2 = 0$ then **cubic (atypical)**. For the semicircle case, $\text{Re } m = \text{const}$ and $f = 1$, so $\tau_2 \neq 0$, thus **quadratic (typical)**.

Precision, rigidity

We prove optimal entry-wise local law

$$\left| G_{jk}(z) - \delta_{jk} m_j(z) \right| \prec \sqrt{\frac{\rho(E)}{N\eta}} + \frac{1}{N\eta}, \quad z = E + i\eta$$

(also for the density and for the "isotropic" version). At the cusps the (current) estimate is slightly weaker than optimal.

In terms of rigidity, i.e. comparing eigenvalues λ_i with the corresponding quantiles γ_i of the limiting density, we have

$$\begin{array}{ll} |\lambda_i - \gamma_i| \prec N^{-1} & \text{bulk} \\ |\lambda_i - \gamma_i| \prec N^{-2/3} & \text{edges (also internal)} \\ |\lambda_i - \gamma_i| \prec N^{-3/5} & \text{cusps} \end{array}$$

Optimal scale at the cusps should be $N^{-3/4}$.

Application: correlated Gaussian matrices

Consider a hermitian matrix with correlated Gaussian entries:

$$\mathbb{E}h_{jk} = 0, \quad \mathbb{E}h_{jk}\overline{h_{j'k'}} = \frac{1}{N} \left(R_{j-j', k-k'} + Q_{j-k', k-j'} \right)$$

where R, Q have a decay

$$\sum_{j,k} \left[|R_{jk}| + |Q_{jk}| \right] (|j| + |k|) < \infty$$

For example, such ensemble can be obtained by filtering

$$H = K \star X + c.c. \quad K_{jk} \text{ kernel with } \sum_{jk} |K_{jk}| \sim 1$$

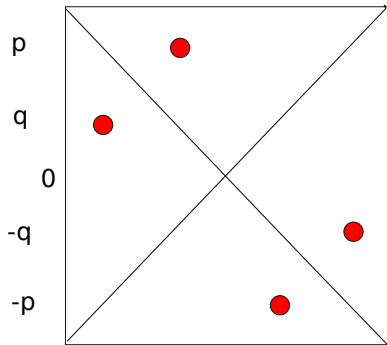
and X has i.i.d. centred Gaussian entries (no symmetry)

$$\mathbb{E}|X_{jk}|^2 = \frac{1}{N}, \quad \mathbb{E}X_{jk}^2 = \frac{\gamma}{N}$$

Then $R = (K \star \tau \bar{K}) + c.c.$ with $(\tau K)_{jk} = K_{-j, -k}$.

Good news: In Fourier space, \widehat{H} has almost independent entries:

$$\widehat{h}_{pq} \perp \widehat{h}_{p'q'} \quad \text{unless} \quad (p', q') \in \{(p, q), (q, p), (-p, -q), (-q, -p)\}$$



Wigner matrix with four-fold symmetry

(hermitian + reflection)

Analysis goes through with this extra symmetry.

Solve the QVE in Fourier space

$$-\frac{1}{m_p(z)} = z + \int_0^1 \widehat{R}_{pq} m_q(z) dq$$

(only R matters, Q is irrelevant)

Apply the previous theorem for $S = \widehat{R}$, get optimal asymptotics for $(\widehat{H} - z)^{-1}$, then Fourier transform back (using the isotropic law).

Theorem [Ajanki-E-Krüger]: Under a nondegeneracy condition, $\inf_p \sup_q |\widehat{R}_{pq}| > 0$, (holds for generic convolution kernels), we have

$$\max_{jk} |G_{jk}(z) - g_{j-k}(z)| \prec \sqrt{\frac{\varrho(E)}{N\eta}}$$

for the resolvent $G = (H - z)^{-1}$ of the correlated Gaussian RM. Here

$$g_k(z) = \int_0^1 e^{-2\pi i k p} m_p(z) dp$$

is the Fourier transform of the solution of the QVE. It inherits the decay of R . Note that G_{jk} is not concentrated to $j = k$.

Similar optimal result for the DOS.

Previous results:

[Schenker, Schulz-Baldes, 2005], [Götze, Naumov, Tikhomirov, 2013] Weak dependence, DOS=sc

[Anderson-Zeitouni, 2008] DOS on macro scale with moment method in case finite range correlation.

[Pastur-Shcherbina, 2011] DOS on macro scale with resolvents.

Local spectral universality

Theorem [Ajanki-E-Krüger-Schnelli] In all models above, bulk local spectral universality holds (in the sense of fixed label or averaged energy).

There is a more general theorem behind which extends previous analysis of the local equilibration of the DBM flow to arbitrary matrix initial condition.

Previous results applied to:

- i) initial matrix follows semicircle [E-Schlein-Yin-Yau]
- ii) deformed Wigner matrices with DOS with a single interval support [Lee-Schnelli-Stetler-Yau]

Summary

- Local laws for Wigner-like matrices (independent entries, arbitrary variance matrix).
- Complete analysis of a NL equation $S_m + z = -\frac{1}{m}$.
- Singularities of the DOS are universal.
- Optimal local laws for the translation invariant correlated Gaussian ensemble.
- Bulk universality in all these models