

Structured Coagulation-Fragmentation Equation in the Space of Radon Measures: Unifying Discrete and Continuous Models*

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Abstract.

We present a structured coagulation-fragmentation model which describes the population dynamics of oceanic phytoplankton. This model is set in the space of Radon measures equipped with the bounded Lipschitz norm and unifies the study of the discrete and continuous coagulation-fragmentation models. We prove that the model is well-posed and show it can reduce down to the classic discrete and continuous coagulation-fragmentation models. To understand the interplay between the physical processes of coagulation and fragmentation and the biological processes of growth, reproduction, and death, we establish a regularity result for the solutions and use it to show that stationary solutions are absolutely continuous under some conditions on model parameters. We present a semi-discrete approximation scheme which conserves mass and use it to present numerical simulations for the model.

Key words. Coagulation-Fragmentation Equations, Structured Populations, Non-negative Radon Measures, Bounded-Lipschitz Norm, Semi-discrete Schemes, Conservation of Mass

AMS subject classifications. 35L60, 35Q92, 92D25

1. Introduction. The discrete coagulation model in the form of a system of differential equations was first introduced by Smoluchowski in his seminal work [54] and was later extended to a continuous setting in the form of an integro-differential equation by Müller [48]. In [12] Blatz and Tobolsky added discrete fragmentation kernels to the literature which were brought into a continuous setting by Melzak [47]. In [5] Ackleh and Fitzpartick introduced the coagulation equations in the context of size-structured population and the fragmentation equation were added to size-structured models by Ackleh in [1]. These models take the form of a nonlinear nonlocal first-order hyperbolic differential equation with a nonlocal boundary condition.

Coagulation-fragmentation equations have been used in many applications in physics, chemistry and biology. In particular, they receive much attention in the study of the population dynamics of phytoplankton [1, 4, 5, 10, 14, 35, 36, 52], which is a vital member of the oceanic ecosystem. Coagulation-fragmentation equations are useful in this application as phytoplankton populations are often modeled as a collection of particles which are held together via an organic glue. Thus, particles can either stick together to form a cell of larger size (coagulate) or fracture off into cells of smaller size (fragment). Coagulation-Fragmentation models are often set with either a continuous size structure [5, 14, 36] or a discrete size structure [9, 15, 42]. In the case of the continuous models, the growth of individual cells through biological means is naturally modeled via a structured partial differential equation. In this work, we extend this idea by presenting a structured coagulation-fragmentation equation in a measure setting with the aim to unify the study of the discrete and continuous equations.

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37 In this paper, we consider the following structured coagulation-fragmentation equation:

$$38 \quad \begin{cases} \partial_t \mu + \partial_x(g(t, \mu)\mu) + d(t, \mu)\mu = K[\mu] + F[\mu], & (t, x) \in (0, T) \times (0, \infty) \\ g(t, \mu)(0)D_{dx}\mu(0) = \int_{\mathbb{R}^+} \beta(t, \mu)(y)\mu(dy), & t \in [0, T]. \\ \mu(0) \in \mathcal{M}^+(\mathbb{R}^+), \end{cases}$$

39 where $\mu(t)$ belongs to $\mathcal{M}^+(\mathbb{R}^+)$, the set of finite nonnegative Borel measures on $\mathbb{R}^+ := [0, +\infty)$.
 40 Here, given a Borel subset $A \subset \mathbb{R}^+$, $\mu_t(A) := \mu(t)(A)$ represents the number of individuals at time
 41 t of structure (e.g. size, age) x in A , and the functions g and d represent the growth and death
 42 rate of individuals at time t of structure x , respectively. Likewise, the function β represents the
 43 reproduction rate of these individuals. More precisely, at time t and distribution $\mu(t)$, an individual
 44 with structure x produces offspring at rate $\beta(t, \mu(t))(x)$. Finally, $D_{dx}\mu(0)$ denotes the Radon-
 45 Nikodym derivative of $\mu(t)$ with respect to the Lebesgue measure, dx , at the point $x = 0$. For
 46 more information about size structured models in a measure setting, we direct the reader to [6, 31].
 47 Finally, K and F are the coagulation and fragmentation terms respectively that we will precisely
 48 define later.

49 The first equation in the above model describes how the number of individuals with structure
 50 x , $\mu(t)(x)$ informally, changes in time t due to the combination of the transport term $\partial_x(g(t, \mu)\mu)$
 51 which moves the distribution μ at velocity g , the death rate which removes individuals from the
 52 system at rate d , the coagulation term $K[\mu]$ which glues individuals together and the fragmentation
 53 term $F[\mu]$ which breaks them. The second equation models the inflow of individuals at the boundary
 54 due to birth. The third equation simply states the regularity of the initial condition.

55 Throughout the literature, there are a variety of assumptions on the coagulation kernel. Com-
 56 mon assumptions include: the kernel being bounded by some combination of linear functions [9, 32];
 57 some ratio of kernel and sizes of particles tending to zero [37, 49]; and, the kernel blows up for small
 58 sizes [18]. Without some additional assumptions on either the kernel or initial condition, the above
 59 assumptions can cause the formation of particles of infinite size. This phenomenon is known as
 60 gelation and has been shown to happen in finite time. Since gelation is not the focus of this paper,
 61 we will require more regularity on our coagulation kernel.

62 Most studies of coagulation-fragmentation equations focus on the case of binary fragmentation;
 63 in other words, when particles only fragment into two smaller units (see [44] and the references
 64 therein, as well as the previously mentioned works). Although the initial work [47] considers the
 65 more general case of multiple fragmentation, where particles can fragment into more than 2 smaller
 66 particles, it is difficult to find many results concerning this case. In the setting of density-based equa-
 67 tions, the authors of [47, 46] work with only an assumption of bounded kernels for both coagulation
 68 and fragmentation. Meanwhile, the work [39] allows for linear growth in the rate of fragmenta-
 69 tion, but requires a bound on the coagulation kernel. The case where both the coagulation and
 70 fragmentation kernels are unbounded is studied in [27, 28].

71 In recent years, the space of Radon measures equipped with the bounded Lipschitz norm has
 72 been used in the study of population dynamics [16, 17, 31, 34]. While many population models
 73 have been studied intensely in this setting, the study of coagulation-fragmentation equations in this
 74 space is sparse. Mild measure solutions to a coagulation-diffusion equation have been obtained in
 75 [49]. The state-space of study was the space of finite measures with absolutely continuous first
 76 marginal and the model does not include any biological processes (i.e. growth, birth, or death).
 77 Existence of solutions to a coagulation-fragmentation equation is obtained in [22] via probabilistic
 78 means. However, authors in [22] only prove existence of a measure solution in the topology of

79 weak convergence and also do not consider any biological processes. The authors in [19] consider
80 a growth-fragmentation equation with a multiple fragmentation kernel identical to that studied
81 in [27]. They cite well-posedness of their model as a consequence of [17] and do not consider a
82 coagulation term. We adopt similar assumptions on our model ingredients, but will prove well-
83 posedness using a fixed-point approach presented in [8]. Finally, for a structured model without
84 coagulation or fragmentation, [34] proves that solutions are absolutely continuous to the left of
85 the zero characteristic curve. Under similar assumptions, we will extend this result to structured
86 coagulation-fragmentation equations.

87 The layout of the paper is as follows. In section 2, we present notation used throughout the
88 paper. In section 3, we reintroduce the model and prove some useful properties of the model
89 ingredients and as well as show the model is indeed well-posed. In section 4, we analyze the
90 interplay between the biological processes (growth, death and birth) and the physical processes
91 (coagulation and fragmentation). In particular, we study their effects on the regularity of solutions
92 to the structured model. In section 5, we show that the classic density and discrete equations are
93 special cases of our model. In section 6, we present a semidiscrete numerical scheme which we
94 test against a few examples providing approximate error in the BL-norm and the numerical order.
95 Finally, in section 7 we will provide discussion of the results and some concluding remarks.

96 **2. Preliminaries and Notation.** In this section, we will provide some preliminary notation. The
97 space of finite Radon measures over $\mathbb{R}^+ := [0, \infty)$ is denoted by $\mathcal{M}(\mathbb{R}^+)$. The non-negative cone
98 of $\mathcal{M}(\mathbb{R}^+)$ will be denoted $\mathcal{M}^+(\mathbb{R}^+)$. Unless otherwise stated, both of these spaces will always be
99 equipped with the Bounded-Lipschitz norm given by

$$100 \quad \|\mu\|_{BL} := \sup_{\|\phi\|_{W^{1,\infty}} \leq 1} \left\{ \int_{\mathbb{R}^+} \phi(x) \mu(dx) : \phi \in W^{1,\infty}(\mathbb{R}^+) \right\}.$$

101 Here, $W^{1,\infty}(\mathbb{R}^+)$ is the usual Sobolev space over \mathbb{R}^+ with codomain \mathbb{R} equipped with the usual
102 norm $\|\phi\|_{W^{1,\infty}} := \|\phi\|_{\infty} + \|\phi'\|_{\infty}$. In the literature, the BL-norm has had a few names such as the
103 flat norm [24, 25], the Dudley norm [21, 23], and the Fortet-Mourier norm [26, 40]. Another norm
104 commonly associated with measures is the total variation norm given by

$$105 \quad \|\nu\|_{TV} = |\nu|(\mathbb{R}^+) = \sup_{\|f\|_{\infty} \leq 1} \left\{ \int_{\mathbb{R}^+} f d\nu : f \in C_c(\mathbb{R}^+) \right\}.$$

106 It should be noted that while over nonnegative measures they are equivalent, the BL-norm and
107 TV-norm are different on the space of signed measures. In particular, for $\mu \in \mathcal{M}(\mathbb{R})$

$$108 \quad \|\mu\|_{BL} \leq \|\mu\|_{TV}.$$

109 We refer the reader to [30] and the references therein for more information.

110 We say a sequence (μ_n) of Radon measures is tight if

$$111 \quad \lim_{x \rightarrow \infty} \sup_n \mu_n([x, \infty)) = 0.$$

112 In $\mathcal{M}^+(\mathbb{R}^+)$, we additionally have that the BL-norm metrizes weak convergence. That is (μ_n)
113 converges weakly to $\mu \in \mathcal{M}^+(\mathbb{R}^+)$ if for every $f \in C_b(\mathbb{R}^+)$,

$$114 \quad \int_{\mathbb{R}^+} f d(\mu_n - \mu) \longrightarrow 0$$

115 as $n \rightarrow \infty$. For more detail, see [31, 29].

116 It is often convenient to use the operator notation in place of integration. That is for a function
117 f , we say

$$118 \quad (\mu, f) := \int_{\mathbb{R}^+} f(x)\mu(dx).$$

119 Finally, we say the flow of a Lipschitz vector field $g(t, x)$ is a function $T_{s,t}^g(x)$ which satisfies

$$120 \quad (3.1) \quad \frac{d}{dt}T_{s,t}^g(x) = g(t, T_{s,t}(x)), \quad T_{s,s}^g(x) = x.$$

121 **3. Structured Coagulation-Fragmentation Equation.** In this section, we establish existence
122 and uniqueness in the space of Radon measures for the structured coagulation-fragmentation equa-
123 tion given by

$$124 \quad (3.1) \quad \begin{cases} \partial_t \mu + \partial_x(g(t, \mu)\mu) + d(t, \mu)\mu = K[\mu] + F[\mu], & (t, x) \in (0, T) \times (0, \infty) \\ g(t, \mu)(0)D_{dx}\mu(0) = \int_{\mathbb{R}^+} \beta(t, \mu)(y)\mu(dy), & t \in [0, T]. \\ \mu(0) \in \mathcal{M}^+(\mathbb{R}^+), \end{cases}$$

125 where

$$126 \quad (3.2) \quad \begin{aligned} \mu &: [0, T] \rightarrow \mathcal{M}^+(\mathbb{R}^+), \\ g, d, \beta &: [0, T] \times \mathcal{M}^+(\mathbb{R}^+) \rightarrow W^{1,\infty}(\mathbb{R}^+), \\ K &: \mathcal{M}^+(\mathbb{R}^+) \rightarrow \mathcal{M}(\mathbb{R}^+), \\ F &: \mathcal{M}^+(\mathbb{R}^+) \rightarrow \mathcal{M}(\mathbb{R}^+). \end{aligned}$$

The model functions g, d , and β are nonnegative and represent the growth, death, and birth functions, respectively. They are assumed to be influenced by both time, t , and the solution to the population model, $\mu(t)$. In applications (e.g., see [2, 3, 17, 20]), it is common to choose β, g and d to depend on a weighted mean of the population in the following form:

$$\beta(t, \mu)(x) = B\left(t, x, \int_{\mathbb{R}^+} K_B(y)d\mu(y)\right)$$

127 and similar expressions for g and d , for given maps $B : [0, T] \times \mathbb{R}^+ \times \mathbb{R}^+ \rightarrow \mathbb{R}^+$ and $K_B : \mathbb{R}^+ \rightarrow \mathbb{R}^+$.
128 Common physically motivated model functions utilize Beverton–Holt type [11] or Ricker type [51]
129 nonlinearities with respect to the weighted mean of the population and of a Von Bertalanffy type
130 [50] model with respect to structure x .

131 The coagulation term is the measure given by

$$132 \quad (3.3) \quad \begin{aligned} K[\mu](\cdot) &= \frac{1}{2} \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa(y, y')\delta_{y+y'}(\cdot)\mu(dy')\mu(dy) - \int_{\mathbb{R}^+} \kappa(y, x)\mu(dy)\mu \\ &=: K^+[\mu] - K^-[\mu], \end{aligned}$$

133 where $\kappa(x, y)$ represents the rate at which individuals of size x coalesce with individuals of size y .
134 The first term in (3.3), K^+ , represents the inflow of individuals due to coagulation. The second
135 term in (3.3), K^- represents the number of individuals lost due to coagulation. Notice that $K^\pm[\mu]$
136 are measures which can be described in a distribution sense by

$$137 \quad (3.4) \quad (K^+[\mu], \phi) = \frac{1}{2} \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa(y, x)\phi(x+y)\mu(dx)\mu(dy).$$

138 and

$$139 \quad (3.5) \quad (K^-[\mu], \phi) = \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa(y, x) \phi(x) \mu(dy) \mu(dx).$$

140 These terms are generalizations of the coagulation terms of the continuous coagulation equation
141 given by

$$142 \quad (3.6) \quad K^+(u)(x) = \frac{1}{2} \int_0^x \kappa(y, x-y) u(x-y) u(y) dy, \quad K^-(u)(x) = u(x) \int_0^\infty \kappa(y, x) u(y) dy.$$

143 Indeed, multiplying $K^+(u)$ by a test function ϕ and integrating we see that

$$144 \quad \frac{1}{2} \int_0^\infty \int_0^x \kappa(y, x-y) u(x-y) u(y) dy \phi(x) dx = \frac{1}{2} \int_0^\infty \int_y^\infty \kappa(y, x-y) \phi(x) u(x-y) dx u(y) dy$$

$$145 \quad = \frac{1}{2} \int_0^\infty \int_0^\infty \kappa(y, x) \phi(x+y) u(x) dx u(y) dy.$$

146 which is $(K^+[\mu], \phi)$ for $\mu = u(y) dy$. An analogous reasoning yields K^- . Notice that if κ is
147 symmetric, i.e. $\kappa(x, y) = \kappa(y, x)$, then

$$148 \quad (3.7) \quad (K[\mu], \phi) = \frac{1}{2} \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa(y, x) [\phi(x+y) - \phi(x) - \phi(y)] \mu(dx) \mu(dy).$$

149 Notice by formally taking $\mu = \sum_{i \in \mathbb{N}} m_i \delta_{x_i}$ we can arrive at the traditional Smoluchowski equations.
150 The fragmentation term is given by

$$151 \quad (3.8) \quad F[\mu](\cdot) = \int_{\mathbb{R}^+} b(y, \cdot) a(y) \mu(dy) - a\mu =: F^+[\mu] - F^-[\mu].$$

Here, $a(y)$ represents the global fragmentation rate of individuals of size y and $b(y, \cdot)$ is a measure supported on $[0, y]$ such that $b(y, A)$ represents the probability a particle of size y fragments to a particle with size in the Borel set A . The positive term, F^+ , represents the inflow of individuals due to fragmentation, and the negative term, F^- , represents the number of individuals lost due to fragmentation. Similar to the coagulation terms, $F^\pm[\mu]$ are measures given explicitly by

$$(F^+[\mu], \phi) = \int_{\mathbb{R}^+} (b(y, \cdot), \phi) a(y) \mu(dy),$$

where $(b(y, \cdot), \phi) = \int_0^y \phi(x) b(y, dx)$, and

$$(F^-[\mu], \phi) = \int_{\mathbb{R}^+} a(y) \phi(y) \mu(dy).$$

152 These terms are a generalization of the multiple fragmentation terms studied in an L^1 setting

$$153 \quad (3.9) \quad F^+(u)(x) = \int_x^\infty b(y, x) a(y) u(y) dy, \quad F^-(u)(x) = a(x) u(x).$$

154 where, following [22], we allow $b(y, \cdot) = b(y, dx)$ to be a non-negative measure supported in $[0, y]$.

155 We impose the following assumptions on the growth, death and birth functions:

156 (A1) For any $R > 0$, there exists $L_R > 0$ such that for all $\|\mu_i\|_{TV} \leq R$ and $t_i \in [0, \infty)$ ($i = 1, 2$)
 157 the following hold

158
$$\|g(t_1, \mu_1) - g(t_2, \mu_2)\|_\infty \leq L_R(|t_1 - t_2| + \|\mu_1 - \mu_2\|_{BL}),$$

159
$$\|d(t_1, \mu_1) - d(t_2, \mu_2)\|_\infty \leq L_R(|t_1 - t_2| + \|\mu_1 - \mu_2\|_{BL}),$$

160
$$\|\beta(t_1, \mu_1) - \beta(t_2, \mu_2)\|_\infty \leq L_R(|t_1 - t_2| + \|\mu_1 - \mu_2\|_{BL}),$$

161 (A2) There exists $\zeta > 0$ such that for all $T > 0$

162
$$\sup_{t \in [0, T]} \sup_{\mu \in \mathcal{M}^+(\mathbb{R}^+)} \|g(t, \mu)\|_{W^{1, \infty}} + \|d(t, \mu)\|_{W^{1, \infty}} + \|\beta(t, \mu)\|_{W^{1, \infty}} < \zeta,$$

163 (A3) For all $(t, \mu) \in [0, \infty) \times \mathcal{M}^+(\mathbb{R}^+)$,

164
$$g(t, \mu)(0) > 0.$$

165 We assume that the coagulation kernel κ satisfies the following assumption:

166 (K) κ is symmetric, nonnegative, bounded by a constant C_κ , and globally Lipschitz with Lip-
 167 schitz constant L_κ .

168 We assume that the fragmentation kernel satisfies the following assumptions:

169 (F1) $a \in W^{1, \infty}(\mathbb{R}^+)$ is non-negative,

170 (F2) for any $y \geq 0$, $b(y, dx)$ is a measure such that

171 (i) $b(y, dx)$ is non-negative and supported in $[0, y]$ so that for all $y > 0$ there exist a
 172 $C_b > 0$ such that $b(y, \mathbb{R}^+) < C_b$,

173 (ii) there exists L_b such that

174
$$\|b(y, \cdot) - b(\bar{y}, \cdot)\|_{BL} \leq L_b|y - \bar{y}|$$

175 (iii) $(b(y, \cdot), x) = y$

176 It follows from (F2) that for any ϕ , $\|\phi\|_{W^{1, \infty}} \leq 1$, the function $\Phi[\phi](y) = (b(y, \cdot), \phi)$ is bounded
 177 Lipschitz with $\|\Phi[\phi](y)\|_{W^{1, \infty}} \leq \bar{C}_b = \max\{C_b, L_b\}$.

178 Given $T \geq 0$, we say a function $\mu \in C([0, T], \mathcal{M}^+(\mathbb{R}^+))$ is a weak solution to (3.1) if for all
 179 $\phi \in (C^1 \cap W^{1, \infty})([0, T] \times \mathbb{R}^+)$, and for all $t \in [0, T]$ the following holds:

180 (3.10)
$$\begin{aligned} & \int_{\mathbb{R}^+} \phi(t, x) \mu_t(dx) - \int_{\mathbb{R}^+} \phi(0, x) \mu_0(dx) = \\ & \int_0^t \int_{\mathbb{R}^+} [\partial_t \phi(s, x) + g(s, \mu_s)(x) \partial_x \phi(s, x) - d(s, \mu_s)(x) \phi(s, x)] \mu_s(dx) ds \\ & + \int_0^t (K[\mu_s] + F[\mu_s], \phi(s, \cdot)) ds + \int_0^t \int_{\mathbb{R}^+} \phi(s, 0) \beta(s, \mu_s)(x) \mu_s(dx) ds. \end{aligned}$$

181 Notice that we can also write model (3.1) with the boundary condition as a source term:

182 (3.11)
$$\partial_t \mu + \partial_x (g(t, \mu) \mu) + d(t, \mu) \mu = K[\mu] + F[\mu] + S(t)[\mu_t]$$

183 where $S(t)[\mu] = \left(\int_0^\infty \beta(t, \mu)(y) \mu(dy) \right) \delta_{x=0}$.

184 The next three propositions discuss useful properties of the source terms.

185 **Proposition 3.1.** For every $\mu \in \mathcal{M}(\mathbb{R}^+)$ we have

$$186 \quad (3.12) \quad \|K[\mu]\|_{TV} \leq \frac{3}{2} C_\kappa \|\mu\|_{TV}^2.$$

187 For every $\mu, \nu \in \mathcal{M}(\mathbb{R}^+)$ with $\|\mu\|_{TV}, \|\nu\|_{TV} \leq R$,

$$188 \quad (3.13) \quad \|K[\mu] - K[\nu]\|_{BL} \leq \bar{L}_{\kappa, R} \|\mu - \nu\|_{BL},$$

189 where $\bar{L}_{\kappa, R}$ is a constant depending only on C_κ, L_κ , and R .

190 *Proof.* To prove (3.12) notice that

$$191 \quad \|K^+[\mu]\|_{TV} \leq \frac{1}{2} \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa(x, y) |\mu|(dx) |\mu|(dy) \leq \frac{1}{2} C_\kappa \|\mu\|_{TV}^2$$

192 and also

$$194 \quad \|K^-[\mu]\|_{TV} \leq \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa(x, y) |\mu|(dx) |\mu|(dy) \leq C_\kappa \|\mu\|_{TV}^2.$$

195 Since $\|K[\mu]\|_{TV} = \|K^+[\mu] - K^-[\mu]\|_{TV} \leq \|K^+[\mu]\|_{TV} + \|K^-[\mu]\|_{TV}$, we obtain (3.12).

196 To prove (3.13), let $\phi \in W^{1, \infty}(\mathbb{R}^+)$ be such that $\|\phi\|_{W^{1, \infty}} \leq 1$. Then

$$198 \quad 2|(K^+[\mu] - K^+[\nu], \phi)| \\ 199 \quad = \left| \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa(y, y') \phi(y + y') \mu(dy) \mu(dy') - \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa(y, y') \phi(y + y') \nu(dy) \nu(dy') \right| \\ 200 \quad = \left| \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa(y, y') \phi(y + y') \mu(dy) (\mu - \nu)(dy') \right. \\ 201 \quad \quad \left. + \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa(y, y') \phi(y + y') \nu(dy') (\mu - \nu)(dy) \right|.$$

202 Since κ is symmetric,

$$204 \quad 2|(K^+[\mu] - K^+[\nu], \phi)| = \left| \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa(y, y') \phi(y + y') (\mu - \nu)(dy) (\mu + \nu)(dy') \right| \\ 205 \quad \leq \int_{\mathbb{R}^+} \left| \int_{\mathbb{R}^+} \kappa(y, y') \phi(y + y') (\mu - \nu)(dy) \right| (|\mu| + |\nu|)(dy').$$

206 For a given $y' \geq 0$, the function $y \mapsto \kappa(y, y') \phi(y + y')$ is bounded Lipschitz with norm $\leq C_\kappa + L_\kappa$.

207 Thus

$$208 \quad 2|(K^+[\mu] - K^+[\nu], \phi)| \leq (C_\kappa + L_\kappa) (\|\mu\|_{TV} + \|\nu\|_{TV}) \|\mu - \nu\|_{BL}.$$

209 Taking the sup over all such ϕ gives

$$\|K^+[\mu] - K^+[\nu]\|_{BL} \leq \frac{1}{2} (C_\kappa + L_\kappa) (\|\mu\|_{TV} + \|\nu\|_{TV}) \|\mu - \nu\|_{BL}.$$

210 In the same way

$$212 \quad |(K^-[\mu] - K^-[\nu], \phi)| = \left| \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa(y, x) \phi(x) \mu(dy) \mu(dx) - \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa(y, x) \phi(x) \nu(dy) \nu(dx) \right| \\ 213 \quad = \left| \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa(y, x) \phi(x) \mu(dy) (\mu - \nu)(dx) + \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa(y, x) \phi(x) (\mu - \nu)(dy) \nu(dx) \right| \\ 214 \quad \leq \int_{\mathbb{R}^+} \left| \int_{\mathbb{R}^+} \kappa(y, x) \phi(x) (\mu - \nu)(dx) \right| |\mu|(dy) + \int_{\mathbb{R}^+} \left| \int_{\mathbb{R}^+} \kappa(y, x) (\mu - \nu)(dy) \right| |\phi(x)| |\nu|(dx) \\ 215 \quad \leq \left((L_\kappa + C_\kappa) \|\mu\|_{TV} + \|\nu\|_{TV} \max\{L_\kappa, C_\kappa\} \right) \|\mu - \nu\|_{BL}$$

217 Combining these two results we see that

$$218 \quad \|K[\mu] - K[\nu]\|_{BL} \leq \bar{L}_{K,R} \|\mu - \nu\|_{BL}. \quad \blacksquare$$

219 Next we have the following proposition concerning the fragmentation term:

220 **Proposition 3.2.** *For any $\mu \in \mathcal{M}(\mathbb{R}^+)$ we have*

$$221 \quad (3.14) \quad \|F[\mu]\|_{TV} \leq (\bar{C}_b + 1) \|a\|_\infty \|\mu\|_{TV}.$$

222 and

$$223 \quad (3.15) \quad \|F[\mu] - F[\nu]\|_{BL} \leq C_{a,b} \|\mu - \nu\|_{BL}.$$

Proof. Clearly,

$$\|F^-[\mu]\|_{TV} \leq \|a\|_\infty \|\mu\|_{TV}$$

and

$$\|F^-[\mu] - F^-[\nu]\|_{BL} \leq \|a\|_{W^{1,\infty}} \|\mu - \nu\|_{BL} = C_a \|\mu - \nu\|_{BL}.$$

Also,

$$\|F^+[\mu]\|_{TV} \leq \|a\|_\infty \|\mu\|_{TV} \|\Phi(1)\|_\infty = \bar{C}_b \|a\|_\infty \|\mu\|_{TV}.$$

and

$$\|F^+[\mu] - F^+[\nu]\|_{BL} \leq \|\mu - \nu\|_{BL} \sup_{\|\phi\|_{W^{1,\infty}} \leq 1} \|\Phi[\phi]a\|_{W^{1,\infty}} = C_{a,b} \|\mu - \nu\|_{BL}.$$

224 The following proposition is immediate from assumptions (A1) and (A2).

225 **Proposition 3.3.** *$S(t)[\mu]$ satisfies the following:*

226• *$S(t)[\mu] \geq 0$ whenever $\mu \geq 0$;*

227• *$\|S(t)[\mu]\|_{TV} \leq \zeta \|\mu\|_{TV}$;*

228• *For any $t \geq 0$ and for any μ, ν with $\|\mu\|_{TV}, \|\nu\|_{TV} \leq R$,*

$$229 \quad \|S(t)[\mu] - S(t)[\nu]\|_{BL} \leq (\zeta + RL_R) \|\mu - \nu\|_{BL}$$

230 **3.1. Well-Posedness of the structured coagulation-fragmentation equation (3.1).** Here, we
231 aim to prove model (3.1) is well-posed. More precisely we prove that

232 **Theorem 3.1.** *Assume that assumptions (A1), (A2), (A3), (K), (F1), (F2) hold. Given an initial
233 condition $\mu_0 \in \mathcal{M}^+(\mathbb{R}^+)$, there exists a unique global solution $\mu \in C([0, \infty), \mathcal{M}^+(\mathbb{R}^+))$ of equation
234 (3.1). Moreover, if μ_0 has finite total mass in the sense that $\int_{\mathbb{R}^+} x \mu_0(dx) < \infty$, then for any $T \geq 0$
235 there exists $C_T > 0$ such that*

$$236 \quad \int_{\mathbb{R}^+} x \mu_t(dx) \leq C_T \quad t \in [0, T].$$

237 *In particular, if $g = d = \beta = 0$ then mass is conserved in the sense that $\int_{\mathbb{R}^+} x \mu_t(dx) = \int_{\mathbb{R}^+} x \mu_0(dx)$
238 for any $t \geq 0$.*

239 *Proof.* Let

$$240 \quad B(t, \mu) := F^+[\mu] + K^+[\mu] + S(t)[\mu]$$

and

$$\bar{N}(t, x, \mu) := -d(t, \mu)(x) - a(x) - \int_{\mathbb{R}^+} \kappa(y, x) \mu(dy).$$

Then equation (3.1) reads

$$\partial_t \mu + \partial_x(g(t, \mu)\mu) = B(t, \mu) + \bar{N}(t, \cdot, \mu)\mu.$$

241 For any $R > 0$, denote $\mathcal{M}_R(\mathbb{R}) := \{\mu \in \mathcal{M}(\mathbb{R}) : \|\mu\|_{TV} \leq R\}$. Notice $\mathcal{M}_R(\mathbb{R})$ is complete
 242 for the BL norm. According to Propositions 3.1, 3.2, and 3.3, $B : \mathbb{R}^+ \times \mathcal{M}(\mathbb{R}) \rightarrow \mathcal{M}(\mathbb{R})$ and
 243 $\bar{N} : \mathbb{R}^+ \times \mathbb{R} \times \mathcal{M}(\mathbb{R}) \rightarrow W^{1,\infty}(\mathbb{R})$ are continuous and satisfy the following properties:

(B1) $B[t, \mu] \in \mathcal{M}^+(R)$ for any $t \geq 0$ if $\mu \in \mathcal{M}^+(\mathbb{R})$,

(B2) for any $R > 0$ there exists $C_{B,R} > 0$ and $L_{B,R} > 0$ such that for any $t \geq 0$ and any $\mu, \tilde{\mu} \in \mathcal{M}_R(\mathbb{R}^d)$,

$$\|B(t, \mu)\|_{TV} \leq C_{B,R}, \quad \text{and} \quad \|B(t, \mu) - B(t, \tilde{\mu})\|_{BL} \leq L_{B,R}\|\mu - \tilde{\mu}\|_{BL}.$$

(N1) for any $R > 0$, there exist $L_{\bar{N},R} > 0$ and $C_{\bar{N},R} > 0$ such that for any $t \geq 0$, $x \in \mathbb{R}$, and any
 $\mu, \tilde{\mu} \in \mathcal{M}_R(\mathbb{R})$,

$$\|\bar{N}(t, \cdot, \mu)\|_{W^{1,\infty}} \leq C_{\bar{N},R} \quad \text{and} \quad |\bar{N}(t, x, \mu) - \bar{N}(t, x, \tilde{\mu})| \leq L_{\bar{N},R}\|\mu - \tilde{\mu}\|_{BL}.$$

245 It follows from standard arguments (e.g. [7, 8] and references therein) that equation (3.1) has a
 246 unique solution $\mu \in C([0, T^*), \mathcal{M}(\mathbb{R}^+))$ which is nonnegative and defined on a maximal time interval
 247 $[0, T^*)$. Moreover, $T^* < \infty$ if and only if $\lim_{t \rightarrow T^*} \|\mu_t\|_{TV} = \infty$. Indeed this follows applying Banach
 248 fixed-point Theorem to the map $\Gamma : X_T \rightarrow X_T$ with

$$249 \quad (3.16) \quad X_T = \{\mu \in C([0, T], \mathcal{M}(\mathbb{R}^+)) : \mu(0) = \mu_0, \|\mu\|_{TV} \leq 2\|\mu_0\|_{TV} \forall t \in [0, T]\},$$

250 and

$$251 \quad (3.17) \quad \Gamma[\mu]_t = T_{0,t}^g \# \mu_0 + \int_0^t T_{s,t}^g \# N(s, \mu) ds,$$

252 where $N(s, \mu) := \bar{N}(s, \cdot, \mu)\mu + B(s, \mu)$, and $T_{s,t}^g$ is the flow of the vector field $(t, x) \rightarrow g(t, \mu_t)(x)$.
 253 We can then prove that taking T small enough, $\Gamma(X_T) \subset X_T$ and Γ is a strict contraction. We then
 254 deduce that (3.1) has a unique solution $\mu \in C([0, T^*), \mathcal{M}(\mathbb{R}^+))$. If moreover $\mu_0 \geq 0$ we can then
 255 prove as [8][Prop. 5.1 and Thm 5.2] that $\mu_t \geq 0$ for any $t < T^*$.

256 Recall that if $T^* < \infty$ then it must be $\lim_{t \rightarrow T^*} \|\mu_t\|_{TV} = \infty$. Thus to prove that $T^* = \infty$, it is
 257 enough to verify that there exists $C > 0$ such that

$$258 \quad (3.18) \quad \|\mu_t\|_{TV} \leq \|\mu_0\|_{TV} \exp(Ct) \quad \text{for any } t \in [0, T^*).$$

259 To begin, we first note for any finite non-negative measure μ ,

$$260 \quad (K[\mu], 1) = -\frac{1}{2} \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa(x, y) \mu(dx) \mu(dy) \leq 0$$

261 and

$$262 \quad (F[\mu], 1) = \int_{\mathbb{R}^+} (b(y, \cdot), 1) a(y) \mu(dy) - \int_{\mathbb{R}^+} a(y) \mu(dy) \leq \int_{\mathbb{R}^+} (C_b - 1) a(y) \mu(dy).$$

263 Therefore, taking $\phi(t, x) \equiv 1$ in (3.10), we can arrive at

$$264 \quad (3.19) \quad \begin{aligned} (\mu_t, 1) &\leq (\mu_0, 1) + \int_0^t \int_{\mathbb{R}^+} [(C_b - 1) a(y) + \beta(s, \mu_s)(y)] \mu_s(dy) ds \\ &\leq (\mu_0, 1) + [(C_b - 1)\|a\|_\infty + \zeta] \int_0^t (\mu_s, 1) ds. \end{aligned}$$

265 The Gronwall inequality then gives (3.18) with $C = (C_b - 1)\|a\|_\infty + \zeta$.

266 Now, assume that $\int_0^\infty x \mu_0(dx) < \infty$. Let $R > 0$ and consider a smooth regularization of the test
 267 function $\phi_R(x) = \min\{x, R\}$ in the weak formulation (3.10). Since $\phi_R(x+y) - \phi_R(x) - \phi_R(y) \leq 0$
 268 for any $x, y \geq 0$, we have from equation (3.7) that $(K[\mu_t], \phi_R) \leq 0$. Moreover, $\phi_R(0) = 0$ and
 269 $\phi_R \geq 0$. We thus obtain

$$270 \quad (\mu_t, \phi_R) \leq (\mu_0, \phi_R) + \int_0^t \int_{\mathbb{R}^+} g(s, \mu_s)(y) \phi_R'(y) \mu_s(dy) ds + \int_0^t \int_{\mathbb{R}^+} (b(y, \cdot), \phi_R) a(y) \mu_s(dy) ds.$$

271 Using (A2) and (3.18), we can bound the 2nd term on the right-hand side by $C_{T,\zeta}$ for $t \in [0, T]$.

272 Using that $\phi_R(x) \leq x$, $(b(y, dx), x) = y$, and (A2), we have

$$273 \quad (\mu_t, \phi_R) \leq (\mu_0, x) + C_{T,\zeta} + \int_0^t \int_{\mathbb{R}^+} ya(y) \mu_s(dy) ds$$

$$274 \quad \leq (\mu_0, x) + C_{T,\zeta} + \|a\|_\infty \int_0^t (\mu_s, x) ds.$$

Passing to the limit $R \rightarrow \infty$ using the Monotone Convergence Theorem, we deduce

$$(\mu_t, x) \leq (\mu_0, x) + C_{T,\zeta} + \|a\|_\infty \int_0^t (\mu_s, x) ds.$$

The Gronwall inequality then gives

$$(\mu_t, x) \leq ((\mu_0, x) + C_{T,\zeta}) e^{\|a\|_\infty t}.$$

275 As a consequence we can use any continuous test-function ϕ with linear growth, i.e. $|\phi(x)| \leq$
 276 $C(1+|x|)$. In particular, we can take $\phi(x) = x$ in equation (3.10). Since $(K[\mu_t], x) = (F[\mu_t], x) = 0$,
 277 we obtain

$$278 \quad (\mu_t, x) = (\mu_0, x) + \int_0^t \int_{\mathbb{R}^+} g(s, \mu_s)(y) \mu_s(dy) ds - \int_0^t \int_{\mathbb{R}^+} xd(s, \mu_s)(x) \mu_s(dx) ds$$

279 In particular, if $g = d = 0$, we have $(\mu_t, x) = (\mu_0, x)$ i.e. mass is conserved for any $t \geq 0$. ■

280 **Remark 3.1.** *In applications the smallest size will not be of size 0 but rather some $x_0 > 0$. Model*
 281 *(3.11) and the Theorem above can be adjusted for such applications by shifting the Dirac measure at*
 282 *0 to x_0 , requiring $g(t, \mu_t)(x_0) > 0$, and requiring $b(y, \cdot)$ to be supported on $[x_0, y]$. In this case, the*
 283 *mass conservation equation would be*

$$284 \quad (\mu_t, x) = (\mu_0, x) + \int_0^t \int_{\mathbb{R}^+} g(s, \mu_s)(y) \mu_s(dy) ds - \int_0^t \int_{\mathbb{R}^+} xd(s, \mu_s)(x) \mu_s(dx) ds$$

$$285 \quad + \int_0^t \int_{\mathbb{R}^+} x_0 \beta(s, \mu_s)(x) \mu_s(dx) ds.$$

286 **3.2. A stability result.** Let us consider a sequence of equations

$$287 \quad (3.20) \quad \begin{cases} \partial_t \mu + \partial_x (g^n(t, \mu) \mu) + d^n(t, \mu) \mu = K^n[\mu] + F^n[\mu], & (t, x) \in (0, \infty) \times (0, \infty) \\ g^n(t, \mu)(0) D_{dx} \mu(0) = \int_{\mathbb{R}^+} \beta^n(t, \mu)(y) \mu(dy), & t \geq 0, \\ \mu^n(0) \in \mathcal{M}^+(\mathbb{R}^+), \quad \int_0^\infty (1+x) \mu^n(0)(dx) < \infty, \end{cases}$$

where

$$K^n[\mu](\cdot) = \frac{1}{2} \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa^n(y, y') \delta_{y+y'}(\cdot) \mu(dy') \mu(dy) - \int_{\mathbb{R}^+} \kappa^n(y, x) \mu(dy) \mu,$$

and

$$F[\mu](\cdot) = \int_{\mathbb{R}^+} b^n(y, \cdot) a^n(y) \mu(dy) - a^n \mu.$$

288 Let us assume that

289 (S1) the functions $g^n, d^n, \beta^n, \kappa^n, a^n, b^n, n \in \mathbb{N}$, satisfy assumptions (A1),(A2),(A3),(K),(F1),(F2),
 290 It then follows from Theorem 3.1 that (3.20) has a unique solution $\mu^n \in C([0, \infty), \mathcal{M}(\mathbb{R}_+))$ such
 291 that $\int_0^\infty x \mu^n(t)(dx) < \infty$. Under some additional assumptions on the coefficients of (3.20) we can
 292 extract from μ^n a subsequence converging to a solution of (3.1).

293 **Theorem 3.2.** *Assume that the functions $g^n, d^n, \beta^n, \kappa^n, a^n, b^n, n \in \mathbb{N}$, satisfy assumptions (S1)*
 294 *and also that*

(S2) *there exists $C > 0$ such that $\|\kappa^n\|_\infty, \|a^n\|_\infty \leq C$ and there exists functions κ, a such that*

$$\kappa^n \rightarrow \kappa, a^n \rightarrow a \quad \text{uniformly on compact sets.}$$

(S3) *there exists $C > 0$ and a function $b : \mathbb{R}_+ \rightarrow \mathcal{M}(\mathbb{R}_+)$ such that $(b^n(y), 1) \leq C$ for any $y \geq 0$ and
 $n \in \mathbb{N}$, and for any $\phi \in C_c^\infty(\mathbb{R}^+)$,*

$$(b^n(y), \phi) \rightarrow (b(y), \phi) \quad \text{uniformly for } y \text{ in a compact set.}$$

(S4) *there exist functions $g, d, \beta : [0, \infty) \times \mathcal{M}^+(\mathbb{R}^+) \rightarrow W^{1,\infty}(\mathbb{R}^+)$ such that for any $t \geq 0$ and any
 sequence of measures $m^n \in \mathcal{M}^+(\mathbb{R}^+)$ converging weakly to $m \in \mathcal{M}^+(\mathbb{R}^+)$ we have*

$$g^n(t, m^n) \rightarrow g(t, m), \quad d^n(t, m^n) \rightarrow d(t, m), \quad \beta^n(t, m^n) \rightarrow \beta(t, m)$$

295 *uniformly on compact sets of \mathbb{R}^+ .*

296 *Concerning the initial condition $\mu^n(0) \in \mathcal{M}^+(\mathbb{R}_+)$, we assume that $\int_{\mathbb{R}^+} (1+x) \mu_0^n(dx) \leq C$ and
 297 $\mu_0^n \rightarrow \mu_0$ in the BL norm for some $\mu_0 \in \mathcal{M}^+(\mathbb{R}^+)$.*

298 *Denote μ^n the solution of (3.20). Then, there exists $\mu \in C(\mathbb{R}^+, \mathcal{M}^+(\mathbb{R}^+))$ such that, along a
 299 subsequence, $\mu^n \rightarrow \mu$ in $C([0, T], \mathcal{M}^+(\mathbb{R}^+))$ for any $T > 0$, and μ is a solution of (3.1).*

Proof. We have

$$(K[\mu^n], 1) = -\frac{1}{2} \int_0^\infty \int_0^\infty \kappa^n(x, y) \mu_t^n(dx) \mu_t^n(dy) \leq 0$$

and

$$|(F^n[\mu_t^n], 1)| \leq \int (|(b^n(y), 1)| + 1) |a^n(y)| d\mu_t^n \leq \sup_{y,n} (\|a^n\|_\infty + |(b^n(y), 1)|) (\mu_t^n, 1) = C(\mu_t^n, 1).$$

Moreover, $(\mu_0^n, 1) \rightarrow (\mu_0, 1)$ so that $(\mu_0^n, 1) \leq C$. Taking $\phi = 1$ in the weak formulation of (3.20) we thus obtain

$$(\mu_t^n, 1) \leq (\mu_0^n, 1) + C \int_0^t (\mu_s^n, 1) ds \leq C + C \int_0^t (\mu_s^n, 1) ds.$$

300 It then follows from Gronwall inequality that for any $T > 0$,

$$301 \quad (3.21) \quad (\mu_t^n, 1) \leq C_T \quad t \in [0, T].$$

302 As in the proof of Theorem 3.1, using $\phi_R(x) = \min\{x, R\}$, $R > 0$, as a test-function we obtain

$$\begin{aligned}
303 \quad (\mu_t^n, \phi_R) &\leq (\mu_0^n, \phi_R) + \int_0^t \int_{\mathbb{R}^+} g^n(s, \mu_s^n)(y) \phi_R'(y) \mu_s^n(dy) ds + \int_0^t \int_{\mathbb{R}^+} (b^n(y, \cdot), \phi_R) a^n(y) \mu_s^n(dy) ds \\
304 \quad &\leq C_T + C \int_0^t (\mu_s^n, x) ds.
\end{aligned}$$

305 Letting $R \rightarrow \infty$ using the monotone convergence Theorem, and then applying Gronwall inequality
306 we obtain

$$307 \quad (\mu_t^n, x) \leq C_T.$$

308 In particular, it follows that $(\mu_t^n)_n$ is tight for any $t \in [0, T]$. Moreover for $0 \leq s < t \leq T$, and any
309 $\phi \in W^{1, \infty}$, $\|\phi\|_{W^{1, \infty}} \leq 1$, we have using (3.21) that

$$\begin{aligned}
310 \quad (\mu_t^n - \mu_s^n, \phi) &= \int_s^t (\mu_\tau^n, g(\tau, \mu_\tau^n) \phi') d\tau - \int_s^t (\mu_\tau^n, d(\tau, \mu_\tau^n) \phi) d\tau + \int_s^t (\mu_\tau^n, \beta(\tau, \mu_\tau^n)) \phi(0) d\tau \\
311 \quad &\quad + \int_s^t (K[\mu_\tau^n], \phi) + (F[\mu_\tau^n], \phi) d\tau \\
312 \quad &\leq 3\zeta C_T(t-s) + \int_s^t 3\|k^n\|_\infty \|\phi\|_\infty + C\|\phi\|_\infty d\tau \leq \bar{C}_T(t-s).
\end{aligned}$$

313 Thus, $\|\mu_t^n - \mu_s^n\|_{BL} \leq \bar{C}_T(t-s)$ so that the sequence $(\mu^n)_n \subset C([0, T], \mathcal{M}(\mathbb{R}^+))$ is uniformly equicon-
314 tinuous. By the Arzela-Ascoli Theorem, for any $T > 0$, we therefore have a convergent subsequence
315 (not relabeled) of the μ_t^n in $C([0, T], \mathcal{M}^+(\mathbb{R}^+))$ which converges to some $\mu \in C([0, T], \mathcal{M}^+(\mathbb{R}^+))$. A
316 diagonal argument gives that $\mu^n \rightarrow \mu$ in $C([0, T], \mathcal{M}^+(\mathbb{R}^+))$ for any $T > 0$.

Since ϕ_R is bounded Lipschitz, we can pass to the limit $n \rightarrow \infty$ in $(\mu_t^n, \phi_R) \leq (\mu_t^n, x) \leq C_T$ to
obtain $(\mu_t, \phi_R) \leq C_T$. Sending $R \rightarrow \infty$ gives that for any $T > 0$,

$$(\mu_t, x) \leq C_T \quad \text{for any } t \in [0, T].$$

317 We now want to pass to the limit $n \rightarrow \infty$ in the equation satisfied by μ^n , namely

$$\begin{aligned}
&\int_{\mathbb{R}^+} \phi(t, x) \mu_t^n(dx) - \int_{\mathbb{R}^+} \phi(0, x) \mu_0^n(dx) = \\
318 \quad (3.22) \quad &\int_0^t \int_{\mathbb{R}^+} [\partial_t \phi(s, x) + g^n(s, \mu_s^n)(x) \partial_x \phi(s, x) - d^n(s, \mu_s^n)(x) \phi(s, x)] \mu_s^n(dx) ds \\
&\quad + \int_0^t (K^n[\mu_s^n] + F^n[\mu_s^n], \phi(s, \cdot)) ds + \int_0^t \int_{\mathbb{R}^+} \phi(s, 0) \beta^n(s, \mu_s^n)(x) \mu_s^n(dx) ds.
\end{aligned}$$

319 Let $\phi \in C_c(\mathbb{R}^+ \times \mathbb{R}^+)$. We pass to the limit in the right-hand side using that $\mu_t^n \rightarrow \mu_t$ for any $t \geq 0$.
320 Since $k^n \rightarrow k$ uniformly on compact sets, $(\mu_s^n, 1) \leq C_T$, and $\mu_s^n \otimes \mu_s^n \rightarrow \mu_s \otimes \mu_s$ weakly, we can pass
321 to the limit

$$\begin{aligned}
322 \quad 2(K[\mu_s^n], \phi) &= \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} (k^n(x, y) - \kappa(x, y)) (\phi(x+y) - \phi(x) - \phi(y)) \mu_s^n(dx) \mu_s^n(dy) \\
323 \quad &\quad + \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa(x, y) (\phi(x+y) - \phi(x) - \phi(y)) \mu_s^n(dx) \mu_s^n(dy) \\
324 \quad &\rightarrow \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa(x, y) (\phi(x+y) - \phi(x) - \phi(y)) \mu_s(dx) d\mu_s(dy) = 2(K[\mu_s], \phi).
\end{aligned}$$

Since $|(K[\mu_s^n], \phi)| \leq C$, we obtain by dominated convergence that

$$\int_0^t (K[\mu_s^n], \phi) ds \rightarrow \int_0^t (K[\mu_s], \phi) ds.$$

325 Similarly, we can pass to the limit in $(F[\mu_s^n], \phi)$ in the same way. Finally, in view of (S4), (3.21)
326 and since ϕ has compact support we have for any $s \geq 0$ that

$$\begin{aligned} 327 & \int_{\mathbb{R}^+} d^n(s, \mu_s^n)(x) \phi(s, x) \mu_s^n(dx) \\ 328 &= \int_{\mathbb{R}^+} [d^n(s, \mu_s^n)(x) - d(s, \mu_s)(x)] \phi(s, x) \mu_s^n(dx) + \int_{\mathbb{R}^+} d(s, \mu_s)(x) \phi(s, x) \mu_s^n(dx) \\ 329 &\rightarrow \int_{\mathbb{R}^+} d(s, \mu_s)(x) \phi(s, x) \mu_s(dx). \\ 330 \end{aligned}$$

Since moreover

$$\left| \int_{\mathbb{R}^+} d^n(s, \mu_s^n)(x) \phi(s, x) \mu_s^n(dx) \right| \leq \zeta \|\phi\|_\infty (\mu_s^n, 1) \leq C_T$$

we obtain by the Dominated Convergence Theorem that

$$\int_0^t \int_{\mathbb{R}^+} d^n(s, \mu_s^n)(x) \phi(s, x) \mu_s^n(dx) ds \rightarrow \int_0^t \int_{\mathbb{R}^+} d(s, \mu_s)(x) \phi(s, x) \mu_s(dx) ds.$$

331 We treat the terms with g^n and β^n in the same way. ■

332 **4. Interplay of Growth, Coagulation, and Fragmentation.** In the recent paper [34], it was
333 shown that the steady state solution of a size-structured population model (i.e. model (3.1) with
334 $K \equiv F \equiv 0$) with positive model ingredients is absolutely continuous with respect to the Lebesgue
335 measure. This leads naturally to studying the effect the physical processes of coagulation and
336 fragmentation would have on such regularity. With this in mind, we present the following theorem:

337 **Theorem 4.1.** *Assume (A1)-(A3), (K), (F1), (F2), and (B2) hold with $g(t, \mu_t) \in C^1(\mathbb{R}^+)$ tak-*
338 *ing strictly positive values, and let μ_t be the solution to (3.1) for some some initial condition μ_0 .*
339 *Moreover, assume each measure $b(y, \cdot)$, $y \geq 0$, is absolutely continuous w.r.t. Lebesgue measure with*
340 *density $b(y, x)$, and that the family $\{b(y, \cdot) : y \geq 0\}$ is uniformly equi-integrable in the sense that*
341 *for any $\varepsilon > 0$, there exists $\delta > 0$ such that for any $V \subset \mathbb{R}^+$ measurable with $|V| < \delta$, there holds*
342 *$b(y, V) = \int_V b(y, x) dx < \varepsilon$. Denote $l_0(t)$ the solution to*

$$343 \begin{cases} \frac{d}{dt} l_0(t) = g(t, \mu(t))(l_0(t)), \\ l_0(0) = 0. \end{cases}$$

344 *Then for any $t > 0$, μ_t is absolutely continuous on $[0, l_0(t))$ with respect to the Lebesgue measure*
345 *(i.e. $\mu_t \ll dx$).*

For simplicity of notation, we will denote

$$\tilde{g}(t, x) := g(t, \mu_t)(x), \quad \tilde{\beta}(s) := \int_0^\infty \beta(s, \mu_s)(y) \mu_s(dy), \quad T_{s,t} := T_{s,t}^{\tilde{g}}.$$

We also recall from equation (3.11) that

$$S(s)[\mu_s] = \tilde{\beta}(s) \delta_{x=0}.$$

346 Before we can prove Theorem 4.1, we first need the following useful lemma:

347 **Lemma 4.1.** *Since $\tilde{g} > 0$, the map $\Phi : s \mapsto T_{s,t}(0)$ is a bijection from $[0, t]$ to $[0, l_0(t)]$. Moreover*

$$348 \quad (4.1) \quad \Phi'(s) = -\tilde{g}(s, 0) \exp \left\{ \int_s^t \partial_x \tilde{g}(\tau, T_{s,\tau}(0)) d\tau \right\} \quad \forall s \in [0, t].$$

349 *Moreover for any $0 < s \leq t$, $T_{s,t} : [0, l_0(s)] \rightarrow [0, l_0(t)]$ is a bijection with*

$$350 \quad (4.2) \quad \frac{d}{dx} T_{s,t}(x) = \exp \left\{ \int_s^t \partial_x \tilde{g}(\tau, T_{s,\tau}(x)) d\tau \right\}.$$

Proof. The bijection property of Φ follows from the uniqueness of trajectories and the definition of $l_0(t)$. As for (4.1), taking the derivative with respect to s in $\frac{d}{dt} T_{s,t}(0) = \tilde{g}(t, T_{s,t}(0))$ yields

$$\frac{d}{dt} \left(\frac{d}{ds} T_{s,t}(0) \right) = \partial_x \tilde{g}(t, T_{s,t}(0)) \frac{d}{ds} T_{s,t}(0).$$

Since \tilde{g} is C^1 in x ,

$$\frac{d}{ds} T_{s,t}(0) = \frac{d}{ds} T_{s,t}(0)|_{t=s} \exp \left\{ \int_s^t \partial_x \tilde{g}(\tau, T_{s,\tau}(0)) d\tau \right\}.$$

351 Since $T_{s,t}(0) = \int_s^t \tilde{g}(\tau, T_{s,\tau}(0)) d\tau$ we have $\frac{d}{ds} T_{s,t}(0)|_{t=s} = -\tilde{g}(s, 0)$ and so we deduce (4.1).

352 The proof of (4.2) is identical, but with taking the derivative with respect to x in $\frac{d}{dt} T_{s,t}(x) =$
353 $\tilde{g}(t, T_{s,t}(x))$ and using that $\frac{d}{dx} T_{s,s}(x) = 1$. ■

354 In particular for any bounded measurable function $\phi : [0, \infty) \rightarrow \mathbb{R}$,

$$355 \quad (4.3) \quad \int_0^t (T_{s,t} \# S(s)[\mu_s], \phi) ds = \int_0^t \tilde{\beta}(s) \phi(\Phi(s)) ds \\ = \int_0^{l_0(t)} \phi(x) \frac{\tilde{\beta}(\Phi^{-1}(x))}{\tilde{g}(\Phi^{-1}(x), 0)} \exp \left\{ - \int_{\Phi^{-1}(x)}^t \partial_x \tilde{g}(\tau, T_{\Phi^{-1}(x),\tau}(0)) d\tau \right\} dx,$$

356 so that $\int_0^t T_{s,t} \# S(s)[\mu_s] ds = \int_0^t T_{s,t} \# \tilde{\beta}(s) \delta_0 ds$ is the function

$$357 \quad (4.4) \quad x \mapsto 1_{[0, l_0(t)]}(x) \frac{\tilde{\beta}(\Phi^{-1}(x))}{\tilde{g}(\Phi^{-1}(x), 0)} \exp \left\{ - \int_{\Phi^{-1}(x)}^t \partial_x \tilde{g}(\tau, T_{\Phi^{-1}(x),\tau}(0)) d\tau \right\}.$$

358 We can now prove Theorem 4.1. The proof we propose is inspired by [55][Lemma 3.5] and
359 [41][Lemma 2.6]. However the presence of the growth term adds new difficulties.

360 *Proof.* Recall that the solution μ was obtained as a fixed point of the map Γ defined in (3.17)
361 namely

$$362 \quad \mu_t = T_t \# \mu_0 + \int_0^t T_{s,t} \# (F^+[\mu_s] + \tilde{\beta}(s) \delta_0) ds + \int_0^t T_{s,t} \# (K^+[\mu_s] - \tilde{A}(s, \cdot) \mu_s) ds$$

where $T_{s,t}$ is the flow of the vector field $(t, x) \rightarrow \tilde{g}(t, x) := g(t, \mu_t)(x)$, and

$$\tilde{A}(t, x) = d(t, \mu_t)(x) + a(x) + \int_{\mathbb{R}^+} \kappa(x, y) \mu_t(dy) \geq 0.$$

363 Notice due to the positivity of the model functions

$$364 \quad (4.5) \quad \mu_t \leq T_t \# \mu_0 + \int_0^t T_{s,t} \# (F^+[\mu_s] + \tilde{\beta}(s) \delta_0) ds + \int_0^t T_{s,t} \# K^+[\mu_s] ds.$$

365 Given some $\delta > 0$ and $s \in [0, t]$, let \mathcal{A}_s be the family of subsets of $[0, l_0(s))$ of the form

$$366 \quad (4.6) \quad A = T_{s, s_1}^{-1}(\cdots (T_{s_{n-1}, s_n}^{-1}(T_{s_n, t}^{-1}(E) - x_n) - x_{n-1}) \cdots) - x_1$$

where $n \in \mathbb{N}_0$, $s \leq s_1 \leq \cdots \leq s_n \leq t$, $x_1, \dots, x_n \geq 0$, and $E \subset [0, l_0(t))$ is a Borel subset with $|E| < \delta$. It is implicitly understood that at each step of the construction of A we take the intersection with $[0, \infty)$. Define then

$$\mathcal{E}(s) := \sup \left\{ \mu_s(A) : A \in \mathcal{A}_s \right\},$$

367 where we extend μ_s to $(-\infty, 0)$ by 0.

368 Notice that $T_s \# \mu_0$ is supported in $[l_0(s), \infty)$ and that any $A \in \mathcal{A}_s$ is a subset of $[0, l_0(s))$. It
369 follows that for any $A \in \mathcal{A}_s$ of the form (4.6) we have by (4.5) that

$$370 \quad (4.7) \quad \mu_s(A) \leq \int_0^s (F^+[\mu_\tau] + \tilde{\beta}(\tau)\delta_0)(T_{\tau, s}^{-1}(A)) d\tau + \int_0^s K^+[\mu_\tau](T_{\tau, s}^{-1}(A)) d\tau.$$

371 For any $0 \leq a \leq b \leq T$ and any subset $B \subset [0, \infty)$ we have by (4.2) and assumption (A2) that

$$372 \quad |T_{a, b}^{-1}(B)| = \int_{\mathbb{R}^+} 1_B(T_{a, b}(y)) dy = \int_{\mathbb{R}^+} 1_B(x) \left| \frac{d}{dx} T_{a, b}^{-1}(x) \right| dx \leq e^{\zeta(b-a)} |B|.$$

Using the translation invariance of Lebesgue measure we then have that the measure of A given by (4.6) can be bounded by

$$|A| \leq e^{\zeta((t-s_n)+(s_n-s_{n-1})+\cdots+(s_1-s))} |E| \leq e^{\zeta(t-s)} \delta \leq C_T \delta.$$

Here and in the sequel of the proof, we denote by C_T any constant depending only on T and the constants appearing in assumptions (A1),(A2),(A3),(K),(F1),(F2). It then follows from (4.4), (A2), (A3) that

$$\int_0^s \tilde{\beta}(\tau)\delta_0(T_{\tau, s}^{-1}(A)) d\tau \leq C_T \delta.$$

373 Moreover

$$374 \quad F^+[\mu_\tau](A) = \int_{\mathbb{R}^+} b(y)(A) a(y) \mu_\tau(dy) \leq \|a\|_\infty \|\mu_\tau\|_{TV} \sup_{y \geq 0} b(y)(A).$$

Since $\|\mu_\tau\|_{TV} \leq C_T$, $\tau \in [0, s]$, we obtain

$$\int_0^s F^+[\mu_\tau](T_{\tau, s}^{-1}(A)) d\tau \leq C_T \sup_{y \geq 0, |A| \leq C_T \delta} b(y)(A).$$

375 If we assume that the family $\{b(y, \cdot)\}_{y \geq 0}$ is uniformly equi-integrable then $\sup_{y \geq 0, |A| \leq C_T \delta} b(y)(A)$
376 goes to 0 as $\delta \rightarrow 0$. We denote $o(1)$ any quantity going to 0 as $\delta \rightarrow 0$ uniformly in $t \in [0, T]$ and A .
377 Coming back to (4.7) we thus obtained so far that

$$378 \quad (4.8) \quad \mu_s(A) \leq o(1) + \int_0^s K^+[\mu_\tau](T_{\tau, s}^{-1}(A)) d\tau. \quad \blacksquare$$

To bound the coagulation term in the right-hand side recall the definition of K^+ :

$$2K^+[\mu_\tau](T_{\tau, s}^{-1}(A)) = \int_{\mathbb{R}^+} 1_{T_{\tau, s}^{-1}(A)}(z+y) \kappa(x, y) \mu_\tau(dz) \mu_\tau(dy) \leq \|\kappa\|_\infty \int_{\mathbb{R}^+} \mu_\tau(T_{\tau, s}^{-1}(A) - y) \mu_\tau(dy).$$

Since $T_{\tau,s}^{-1}(A) - y \in \mathcal{A}_\tau$ we obtain

$$2K^+[\mu_\tau](T_{s,t}^{-1}(A)) \leq \|\kappa\|_\infty \mathcal{E}(\tau) \int_{\mathbb{R}^+} \mu_\tau(dy)$$

so that

$$K^+[\mu_\tau](T_{\tau,s}^{-1}(A)) \leq C_T \mathcal{E}(\tau).$$

Coming back to (4.8) we obtain

$$\mu_s(A) \leq o(1) + C_T \int_0^s \mathcal{E}(\tau) d\tau.$$

Since this holds for any $A \in \mathcal{A}_s$ and any $s \leq t$ we deduce

$$\mathcal{E}(t) \leq o(1) + C_T \int_0^t \mathcal{E}(\tau) d\tau$$

which yields by Gronwall inequality

$$\mathcal{E}(t) = o(1).$$

In particular, since $E \in \mathcal{A}_t$,

$$\mu_t(E) = o(1) \quad \forall E \subset [0, l_0(t)), |E| < \delta.$$

379 It follows that μ_t is absolutely continuous on $[0, l_0(t))$ for any $t > 0$.

380 This leads us to the following corollary about the regularity of a steady state solution to model
381 3.1.

382 **Corollary 4.1.** *Let the assumptions of Theorem 4.1 hold with g, d, β dependent on time only*
383 *through μ_t (i.e. $g(t, \mu_t) = g(\mu_t)$ etc.) and assume $\mu \in \mathcal{M}^+(\mathbb{R}^+)$ be a steady state solution of*
384 *model 3.1. Then μ is absolutely continuous with respect to the Lebesgue measure. Furthermore, μ*
385 *satisfies*

$$386 \int_{\mathbb{R}^+} g(\mu)(x) \mu(dx) = \int_{\mathbb{R}^+} x d(\mu)(x) \mu(dx).$$

387 *Proof.* The proof follows from similar arguments of Proposition 2.6 in [34] with making use of
388 $g(\mu)(x) > 0$ for all x . Indeed, since $g(\mu)(x) > 0$ for all x we have

$$389 \lim_{t \rightarrow \infty} l^0(t) = \infty.$$

390 Theorem 4.1 then implies a solution μ_t is absolutely continuous on the interval $[0, l^0(t))$. Thus, the
391 steady state solution $\mu_t = \mu$ is absolutely continuous on $[0, \infty)$. The mass conservation equation
392 follows from Theorem 3.1. ■

393 **5. From Measure Equation to Discrete and Continuous Equations.** It is often claimed that
394 one of the many benefits of population models set in measure spaces is the unification of the study
395 of discrete and continuous structure. In this section, we demonstrate this property by showing that
396 model (3.1) includes as special cases the discrete Smolukowski equations [54] and the continuous
397 Müller model [48].

398 **5.1. Continuous Density Model.** In this subsection, we briefly recover the continuous density
 399 equation studied in [1, 5, 14, 48] from (3.1). This follows naturally under the following assumptions:

- 400 (B1) μ_0 is absolutely continuous with respect to the Lesbesgue measure,
 401 (B2) $b(y, \cdot)$ is absolutely continuous with respect to the Lesbesgue measure.

402 Then by undoing the derivations of (3.3) and (3.8), one arrives at the density equations (3.6) and
 403 (3.9) covered in the aforementioned works. In particular, we can recover the binary fragmentation
 404 kernels studied in [1, 14, 36] by taking

$$405 \quad (5.1) \quad a(y) = \frac{1}{2} \int_0^y \gamma(y-s, s) ds, \quad b(y, \cdot) = \frac{\gamma(x, y-x)}{a(y)} dx$$

406 where the function $\gamma(x, y)$ models the rate at which a particles of size $x + y$ fragment into particles
 407 of size x and y .

408 **5.2. Discrete Equation.** In this subsection, we show under certain assumptions, model (3.1)
 409 will reduce to the discrete coagulation-fragmentation equation discussed in [9, 54]. To obtain these
 410 equations, we set $g(t, \mu) = \beta(t, \mu) \equiv 0$ for the remainder of this section. To this end, suppose that
 411 the measures μ_0 and $b(y, \cdot)$ are supported on $h\mathbb{N} = \{h, 2h, \dots\}$ for some fixed $h > 0$ i.e.

$$412 \quad (C1) \quad \mu_0 = \sum_{i \in \mathbb{N}} m_i^0 \delta_{ih} \text{ where for each } i, m_i^0 \in \mathbb{R}^+,$$

$$413 \quad (C2) \quad b(y, \cdot) = \sum_{i \in \mathbb{N}} b_i(y) \delta_{ih}.$$

414 We then have the following result:

415 **Theorem 5.1.** *Let assumptions (A1), (A2), (K), (F1), (F2), (C1), (C2), and (C3) hold. Then for*
 416 *any $t \in [0, \infty)$, the solution μ_t of (3.1) is supported on $h\mathbb{N}_0$:*

$$417 \quad (5.2) \quad \mu_t = \sum_{l \in \mathbb{N}} m_l(t) \delta_{lh},$$

418 where the $m_l(t)$, $l \in \mathbb{N}$, satisfy the discrete coagulation-fragmentation equation

$$419 \quad (5.3) \quad \begin{aligned} & \frac{d}{dt} m_l(t) + d(t, \mu_t)(lh) m_l(t) \\ &= \frac{1}{2} \sum_{i=1}^{l-1} m_i(t) m_{l-i}(t) \kappa(ih, (l-i)h) - \sum_{i=1}^{\infty} \kappa(ih, lh) m_i(t) m_l(t) \\ & \quad + \sum_{i \geq l} b_l(ih) a(ih) m_i(t) - a(lh) m_l(t) \end{aligned}$$

420 with initial condition $m_l(0) = m_l^0$.

Proof. It is clear from Theorem 3.1 that (3.1) has a unique solution $\mu \in C([0, \infty), \mathcal{M}^+(\mathbb{R}^+))$.
 Moreover, according to the proof of Theorem 3.1, μ is a fixed-point of Γ defined in (3.17). Since
 $g = 0$, $T_{s,t}^g$ is the identity map. Thus Γ is simply given by

$$\Gamma[\nu]_t = \mu_0 + \int_0^t \{F[\nu_s] + K[\nu_s] + S(s)[\nu_s] - d(s, \nu_s)\nu_s\} ds$$

421 for any $\nu \in C([0, \infty), \mathcal{M}(\mathbb{R}^+))$. Notice that if ν_t is supported in $h\mathbb{N}$ for any s then so is $\Gamma[\nu]_t$
 422 (concerning K^+ notice this follows from the fact that $h\mathbb{N} + h\mathbb{N} \subset h\mathbb{N}$). We can thus replace X_T in
 423 (3.16) by

$$424 \quad (5.4) \quad X_T = \{\mu \in C([0, T], \mathcal{M}(h\mathbb{N})) : \mu(0) = \mu_0, \|\mu\|_{TV} \leq 2\|\mu_0\|_{TV} \forall t \in [0, T]\},$$

425 and repeat the proof of Theorem 3.1 verbatim to obtain that μ_t is supported in $h\mathbb{N}$ for any $t \geq 0$.
 426 It follows that μ_t can be written as in (5.2). Equation (5.3) follows from (3.10) taking a C^1 test-
 427 function, ϕ , constant in time and supported in $(lh - h, lh + h)$ such that $\phi(lh) = 1$. ■

428 **6. Numerical Methods and Results.** In this section, we present a semidiscrete scheme for
 429 a coagulation-fragmentation equation based on (5.3) and Theorem 5.1 as well as provide some
 430 numerical results based on this scheme. For the rest of this section, we assume that $\beta(t, \mu) =$
 431 $g(t, \mu) \equiv 0$.

432 **6.1. A semi-discrete numerical scheme.** We consider equation (3.1) with $\int_{\mathbb{R}^+} (1+x)\mu_0(dx) <$
 433 ∞ and we assume that assumptions (A1),(A2),(A3),(K),(F1),(F2) hold. We present a semi-discrete
 434 scheme inspired by [43].

Consider the grid $h\mathbb{N}_0$ and the cell $\Lambda^h(i)$ centered at the grid point ih defined by

$$\Lambda^h(i) := [hi - h/2, hi + h/2), \quad i \geq 1, \quad \Lambda^h(0) = [0, h/2).$$

We define the discretization of the initial condition $\bar{\mu}_0 \in \mathcal{M}^+(\mathbb{R}^+)$ with respect to the grid $h\mathbb{N}_0$ by

$$\mu_0^h = \sum_{i \geq 0} \mu_0^h(i) \delta_{hi}, \quad \mu_0^h(i) = \mu_0(\Lambda^h(i)).$$

435 We want to approximate the solution μ_t of (3.1) by measures μ_t^h supported in $h\mathbb{N}_0$ and solution
 436 of some discretized equation. We first approximate the model coefficients κ , a , b as follow. First we
 437 define

$$438 \quad a_i^h = \frac{1}{h} \int_{\Lambda^h(i)} a(y) dy, \quad \kappa_{i,j}^h = \frac{1}{h^2} \int_{\Lambda^h(i) \times \Lambda^h(j)} \kappa(x, y) dx dy$$

439 for $i, j \geq 1$, and

$$440 \quad a_0^h = \frac{2}{h} \int_{\Lambda^h(0)} a(y) dy, \quad \kappa_{0,0}^h = \frac{4}{h^2} \int_{\Lambda^h(0) \times \Lambda^h(0)} \kappa(x, y) dx dy$$

441 (with the natural modifications for $\kappa_{0,j}^h$ and $\kappa_{i,0}^h$, $i \geq 1$). We then let $a^h \in W^{1,\infty}(\mathbb{R}^+)$ and $\kappa^h \in$
 442 $W^{1,\infty}(\mathbb{R}^+ \times \mathbb{R}^+)$ be the linear interpolation of the a_i^h and $\kappa_{i,j}^h$ respectively. Finally, we define the
 443 measure $b^h(jh, \cdot) \in \mathcal{M}^+(h\mathbb{N})$ by

$$444 \quad b^h(jh, \cdot) = \sum_{i \leq j} b(jh, \Lambda^h(i)) \delta_{ih}$$

and then $b^h(x, \cdot) \in \mathcal{M}^+(h\mathbb{N}_0)$ for $x \geq 0$ as the linear interpolate between the $b^h(jh, \cdot)$. We define
 the corresponding coagulation and fragmentation operators K^h and F^h by

$$(K^h[\mu], \phi) = \frac{1}{2} \int_{\mathbb{R}^+} \int_{\mathbb{R}^+} \kappa^h(y, x) [\phi(x+y) - \phi(x) - \phi(y)] \mu(dx) \mu(dy),$$

$$F^h[\mu](\cdot) = \int_{\mathbb{R}^+} b^h(y, \cdot) a^h(y) \mu(dy) - a^h \mu.$$

Notice that K^h, a^h, b^h satisfy (K),(F1),(F2)(i),(F2)(ii), (C1),(C2),(C3). However (F2)(iii) is
 only satisfied up to an error of order h , namely

$$|(b^h(y, \cdot), x) - y| \leq Ch \quad \text{for any } y \geq 0,$$

445 where the constant C depends only on C_b given by (F2)(i). Indeed recalling that for any $j \geq 0$ the
 446 measure $b(jh, \cdot)$ is non-negative and supported in $[0, jh]$ we have

$$447 \quad |(b^h(jh, \cdot), x) - jh| = |(b^h(jh, \cdot), x) - (b(jh, \cdot), x)| \leq \sum_{i \leq j} \int_{\Lambda^h(i)} |ih - x| b(jh, dx)$$

$$448 \quad \leq \frac{h}{2} b(jh, \mathbb{R}^+) \leq \frac{1}{2} C_b h.$$

450 The result follows recalling that for $y \in [jh, (j+1)h]$ we have $b^h(y, \cdot) = \frac{1}{h} [b^h((j+1)h, \cdot) - b^h(jh, \cdot)](y -$
 451 $jh) + b^h(jh, \cdot)$.

452 It then follow from Theorem 5.1 that (3.1) with $g = d = \beta = 0$, $K = K^h$, $F = F^h$ has a unique
 453 solution $\mu \in C([0, \infty), \mathcal{M}^+(\mathbb{R}^+))$ which is supported on $h\mathbb{N}$:

$$454 \quad (6.1) \quad \mu_t^h = \sum_{l \in \mathbb{N}_0} m_l^h(t) \delta_{lh},$$

455 where the $m_l^h(t)$, $l \in \mathbb{N}_0$, satisfy the discrete coagulation-fragmentation equation

$$456 \quad (6.2) \quad \frac{d}{dt} m_l^h(t) = \frac{1}{2} \sum_{i=1}^{l-1} m_i^h(t) m_{l-i}^h(t) \kappa_{i,l-i}^h - \sum_{i=1}^{\infty} \kappa_{i,l}^h m_i^h(t) m_l^h(t)$$

$$+ \sum_{i \geq l} b(ih, \Lambda^h(l)) a_i^h m_i^h(t) - a_l^h m_l^h(t)$$

457 with initial condition $m_l^h(0) = m_0^h(l)$. Notice that the first two terms on the right hand side of (6.2)
 458 make up the discrete Smoluchowski equations and therefore these terms conserve mass. Indeed,
 459 multiplying by $x_l := lh$ and summing over $l = 1, 2, \dots$ we have

$$460 \quad \frac{1}{2} \sum_{l=1}^{\infty} \sum_{i=1}^{l-1} x_l m_i^h(t) m_{l-i}^h(t) \kappa_{i,l-i}^h - \sum_{l=1}^{\infty} \sum_{i=1}^{\infty} x_l \kappa_{i,l}^h m_i^h(t) m_l^h(t)$$

$$461 \quad (6.3) \quad = \frac{1}{2} \sum_{i=1}^{\infty} \sum_{j=1}^{\infty} (x_i + x_j) m_i^h(t) m_j^h(t) \kappa_{i,j}^h - \sum_{l=1}^{\infty} \sum_{i=1}^{\infty} x_l \kappa_{i,l}^h m_i^h(t) m_l^h(t)$$

$$462 \quad = 0.$$

464 However, since $|(b^h(y, \cdot), x) - y| = O(h)$ it is clear that the fragmentation terms only conserve mass
 465 up to an error of order h .

466 To study the limit of μ_t^h as $h \rightarrow 0$ we first state the following properties:

467 **Proposition 6.1.** *The following holds:*

- 468 (i) $\lim_{h \rightarrow 0} \|\mu_0^h - \mu_0\|_{BL} = 0$ and $\int_{\mathbb{R}^+} (1+x) \mu_0^h(dx) \leq C$,
- 469 (ii) $a^h \rightarrow a$, $\kappa^h \rightarrow \kappa$ uniformly on compact sets, and $a^h, \kappa^h \leq C$.
- 470 (iii) for any $\phi \in W^{1,\infty}(\mathbb{R})$, $(b^h(x), \phi) \rightarrow (b(x), \phi)$ uniformly for x in a compact set.

471 *Proof.* For any $\phi \in W^{1,\infty}(\mathbb{R}^+)$, $\|\phi\|_{W^{1,\infty}} \leq 1$, we have

$$472 \quad (\mu_t^h - \mu_t, \phi) = \sum_{i \geq 0} \int_{\Lambda_i(h)} \phi(ih) - \phi(x) \mu_0(dx) \leq \sum_{i \geq 0} \int_{\Lambda_i(h)} |ih - x| \mu_0(dx) \leq \frac{h}{2} \mu_0(\mathbb{R}^+).$$

Moreover

$$\int_{\mathbb{R}^+} x \mu_0^h(dx) = \sum_{i \geq 0} \int_{\Lambda_i(h)} ih \mu_0(dx) = \sum_{i \geq 0} \int_{\Lambda_i(h)} x \mu_0(dx) + O(h) = (\mu_0, x) + O(h)$$

473 which proves (i).

474 Concerning (ii), since $0 \leq a, \kappa \leq C$, we have $0 \leq a^h, \kappa^h \leq C$. Moreover, let $x \in [nh, mh]$ for
 475 some $n \neq m \in \mathbb{N}_0$. Then letting $\chi_A(x)$ represent the characteristic function over the set A , we have

$$\begin{aligned}
 476 \quad \|a^h - a\|_\infty &\leq \sum_{i=n}^m \left| (a_{i+1}^h - a_i^h) \left(\frac{x - ih}{h} \right) + a_i - a(x) \right| \chi_{[ih, (i+1)h)}(x) \\
 477 \quad &\leq \sum_{i=n}^m |a(x_{i+1}) - a(x_i) + a(x_i) - a(x) + O(h)| \chi_{[ih, (i+1)h)}(x) \\
 478 \quad &\leq \|a\|_{W^{1,\infty}} 2h(m-n) + O(h).
 \end{aligned}$$

480 Finally for (iii) again assume $x \in [nh, mh]$, then for $\phi \in W^{1,\infty}(\mathbb{R})$ we have

$$\begin{aligned}
 481 \quad (b^h(x) - b(x), \phi) &= \sum_{j=n}^m \left[(b_{j+1}^h - b_j^h, \phi) \left(\frac{x - jh}{h} \right) + (b_j^h - b(x), \phi) \right] \chi_{[jh, (j+1)h)}(x) \\
 482 \quad &\leq \sum_{j=n}^m \left[\sum_{i \leq j+1} b((j+1)h, \Lambda^h(i)) \phi(ih) - \sum_{i \leq j} b((j)h, \Lambda^h(i)) \phi(ih) + (b_j^h - b(x), \phi) \right] \\
 483 \quad &= \sum_{j=n}^m [(b((j+1)h) - b(jh), \phi) + (b(jh) - b(x), \phi) + O(h)] \chi_{[jh, (j+1)h)}(x). \\
 484
 \end{aligned}$$

485 Making use of assumption (F2), we have

$$486 \quad (b^h(x) - b(x), \phi) \leq 2L_b \|\phi\|_{W^{1,\infty}} h |m - n|,$$

487 which completes the proof. ■

488 It follows from this proposition that the assumption of Theorem 3.2 are satisfied. Thus, we
 489 deduce that μ^h converges along a subsequence $h \rightarrow 0$ to μ solution of equation (3.1). Since this
 490 equation has a unique solution, the whole sequence μ^h converges to μ :

491 **Theorem 6.1.** *The measure $\mu_t^h = \sum_{i \geq 0} m_i^h(t) \delta_{ih}$ where the m_i^h solve (6.2) converges to the solu-*
 492 *tion μ_t of equation (3.1) in $C([0, T], \mathcal{M}(\mathbb{R}^+))$ for any $T > 0$.*

493 We can thus think of the system (6.2) as a semi-discrete scheme for solving equation (3.1). One
 494 could combine this semidiscrete scheme with any ordinary differential equation scheme (e.g. any
 495 Runge-Kutta Method) to arrive at a fully discrete scheme. Convergence for such a scheme then
 496 follows from a standard triangle inequality argument. In the next section we present some numerical
 497 experiments to evaluate the quality of such a scheme.

498 **Remark 6.1.** *One can easily include the case $\beta, d > 0$ as these terms do not affect the discrete*
 499 *structure of the solution. However, in the case of additionally assuming $g > 0$, it is not true that*
 500 *the solution is discrete for all time. This result was shown for structured population models (without*
 501 *coagulation and fragmentation) in [34] and with coagulation-fragmentation in Section 4.*

502 **6.2. Mass Conserving Fragmentation Term.** To remedy the error generated in mass conser-
 503 vation of the scheme discussed in the previous section, we propose a new approximation of $b(y, dx)$
 504 in the form $b^h(y, \cdot) = \sum_{j=1}^{\infty} \alpha_j(y) \delta_{x_j}$ for which the following holds:

$$505 \quad \sum_{j=1}^{\infty} \alpha_j(y) x_j = (b(y, \cdot), x).$$

506 A natural choice of $\alpha_j(y)$ is given by

$$507 \quad \alpha_j(y) = \frac{1}{x_j} \int_{\Lambda^h(j)} xb(y, dx).$$

508 This approximation results in a mass conserving scheme at the expense of requiring a minimum
509 positive size x_0 . We have the following result:

510 **Proposition 6.2.** *Assume there is a positive minimum size $x_0 > 0$ and therefore the points $x_j =$
511 $x_0 + jh$. Then*

$$512 \quad \|b^h(y, \cdot) - b(y, \cdot)\|_{BL} \longrightarrow 0 \quad \text{as} \quad h \longrightarrow 0.$$

513 *Proof.* Taking $\phi(x) \in W^{1,\infty}(\mathbb{R})$ with $\|\phi\|_{W^{1,\infty}} \leq 1$ and letting $\phi_j := \phi(x_j)$ we have

$$\begin{aligned} 514 \quad (b^h(y, \cdot) - b(y, \cdot), \phi) &= \sum_{i=1}^{\infty} \int_{\Lambda(ih)} \frac{\phi_i}{x_i} x - \phi(x) b(y, dx) \\ 515 &= \sum_{i=1}^{\infty} \int_{\Lambda(ih)} \frac{\phi_i x - \phi(x) x_i}{x_i} b(y, dx) \\ 516 &= \sum_{i=1}^{\infty} \int_{\Lambda(ih)} \frac{\phi_i(x - x_i)}{x_i} + \frac{(\phi_i - \phi(x)) x_i}{x_i} b(y, dx). \end{aligned}$$

518 Since $0 < x_0 \leq x_i$ the first term is bounded and making use of the Lipschitz property of ϕ we have

$$\begin{aligned} 519 \quad \sum_{i=1}^{\infty} \int_{\Lambda(ih)} \frac{\phi_i(x - x_i)}{x_i} + \frac{(\phi_i - \phi(x)) x_i}{x_i} b(y, dx) &\leq \sum_{i=1}^{\infty} \int_{\Lambda(ih)} \left(\frac{\phi_i}{x_0} + 1\right) \frac{h}{2} b(y, dx) \\ 520 &\leq \left(\frac{1}{2x_0} + \frac{1}{2}\right) C_b h \quad \blacksquare \end{aligned}$$

522 Therefore by the same arguments in the section above, we can conclude that a scheme with this
523 term will converge to the solution of equation (3.1) with $g = d = \beta = 0$.

524 The standard kernel taken for a structure domain \mathbb{R}^+ is given by $b(y, dx) = \frac{2}{y} dx$. For the domain
525 $[x_0, \infty)$, an example of a kernel which satisfies assumption (F2) is given by

$$526 \quad (6.4) \quad b(y, dx) := \frac{2q}{y - x_0} \left(\frac{x - x_0}{y - x_0}\right)^{q-1} dx, \quad q = 1 - \frac{2x_0}{y}.$$

527 Notice, that if $x_0 = 0$, then the above kernel reduces to $\frac{2}{y} dx$. It should be noted that it is important
528 to calculate $\alpha_j(y)$ exactly when implementing the scheme. Otherwise, numerical integration error
529 may be introduced resulting in lack of mass conservation.

530 **6.3. Numerical Results.** In this section, we test the semidiscrete scheme against some com-
531 monly used examples. We begin by testing the coagulation and fragmentation portions of the
532 scheme separately. We implement the semidiscrete scheme using MATLAB's ode45 function. In
533 each example, we present the exact solution at time $T = 1$ plotted against the structure variable,
534 x , the absolute value difference of the numerical and exact solution, and the relative mass between
535 the numeric and exact solutions plotted against time. We remark that for examples with only coag-
536 ulation, the semi-discrete scheme (6.2) conserves mass (i.e. (6.3)); therefore, any change of mass is
537 due to simulating infinite domain problems over a finite interval. Where it is applicable, we provide
538 a table calculating the BL-norm and numerical order of the scheme. The BL-norm is approximated
539 by the algorithm provided in [33], while the numerical order of the scheme is calculated using the
540 standard calculation:

$$541 \quad \log_2(\|\mu_t - \mu_t^{2h}\|_{BL} / \|\mu_t - \mu_t^h\|_{BL}).$$

542 **6.3.1. Coagulation and Fragmentation Examples.** In this section we presented several numer-
543 ical example focused on coagulation and fragmentation processes.

544 *Example 1.* For the first example, we take the coagulation kernel $\kappa(x, y) \equiv 1$ with $\mu_0 = e^{-x} dx$
545 and all other model ingredients are set to 0. This problem has an exact solution

$$546 \mu_t = \left(\frac{2}{2+t} \right)^2 \exp \left(-\frac{2}{2+t} x \right) dx$$

547 see [38] for more details. Numerical simulations for this example are presented in Figure 1 with
548 $\Delta x = 1/40$ and the BL error and order of conference are presented in Table 1. Simulation are
549 performed over the finite domain $x \in [0, 20]$.

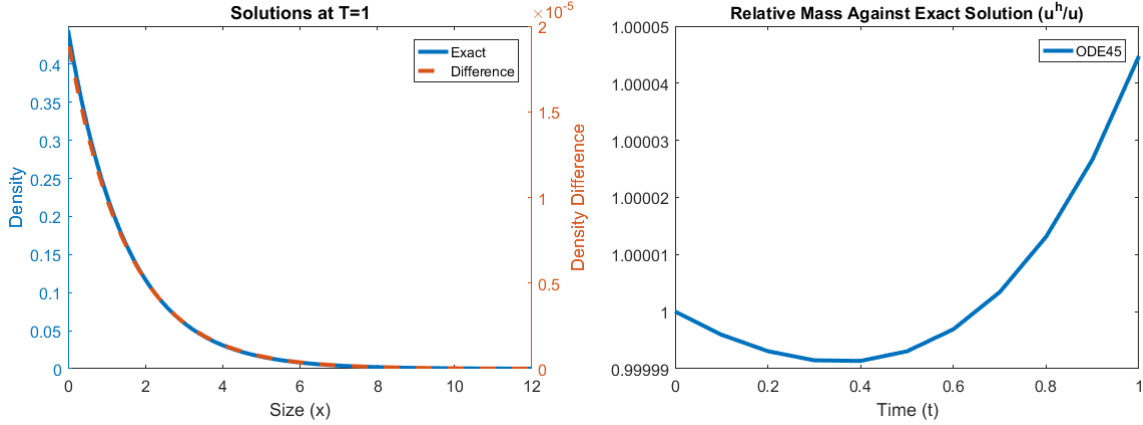


Figure 1: For Example 1 we present on the left side the exact solution (solid line) and the absolute value of the difference between the exact and numerical solution (dashed-line). On the right side we present the relative mass.

Number of Points	BL-Error	Order
40	0.0072641	N/A
80	0.0019723	1.8809
160	0.0005119	1.9459
320	0.00013018	1.9754
640	0.000032716	1.9924
1280	0.0000080986	2.0143

Table 1: Error and numerical order of convergence calculated for Example 1.

550 *Example 2.* Although our theory does not cover the phenomenon of gelation, we include a
551 numerical example showing how the semi discrete scheme handles such kernels. In this example, we
552 take $\kappa(x, y) = xy$ with $\mu_0 = e^{-x}/x dx$. This has exact solution(see e.g. [38].)

$$553 \mu_t = e^{-Tx} \frac{I_1(2xt^{1/2})}{x^2 t^{1/2}} dx,$$

554 where

$$555 T = \begin{cases} 1+t & t \leq 1 \\ 2t^{1/2} & \text{otherwise} \end{cases}$$

556 and

557

$$I_1(x) = \frac{1}{\pi} \int_0^\pi e^{x \cos(\theta)} d\theta.$$

558 Numerical simulations for this example are presented in Figure 6.3.1 with $\Delta x = 1/40$ and the BL
 559 error and order of conference are presented in Table 2. For the order of convergence, the simulations
 560 are performed over the finite domain $x \in [10^{-2}, 20]$.

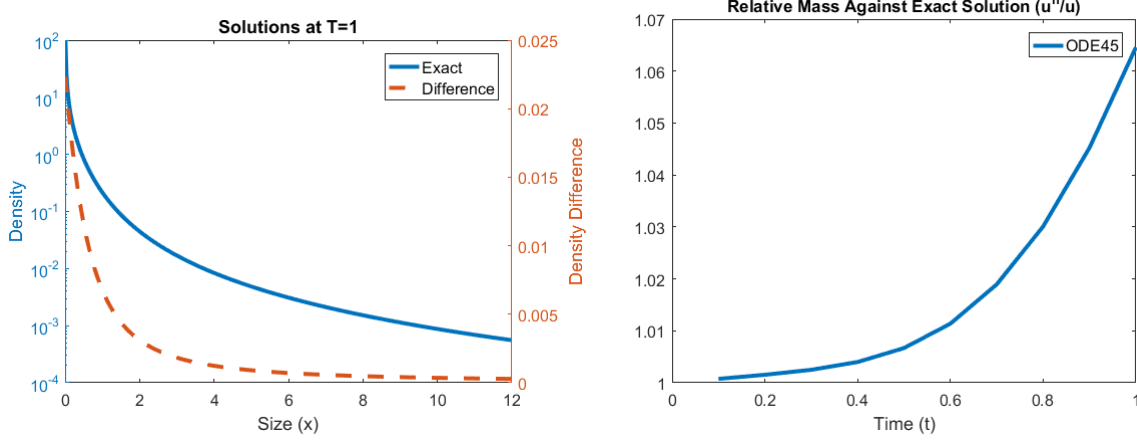


Figure 2: For Example 2 we present on the left side the exact solution (solid line) and the absolute value of the difference between the exact and numerical solution (dashed-line). On the right side we present the relative mass.

561 *Example 3.* In this example we consider fragmentation. We let $b(y, \cdot) = \frac{2}{y} dx$ and $a(x) = x$. As
 562 given in [53], this problem has an exact solution of

563

$$\mu_t = (1 + t)^2 \exp(-x(1 + t)) dx.$$

564 Numerical simulations for this example are presented in Figure 3 with $\Delta x = 1/40$ and the BL error
 565 and order of conference are presented in Table 2. Although convergence for the mass conserving
 566 fragmentation scheme is only shown for positive minimum mass, it still seems to preform well for
 567 the simulations below. Solving the fragmentation terms exactly leads to an $O(h^2)$ term in the last
 568 subinterval (where $y = x_j := j\Delta x$). Explicitly, we have

569

$$\alpha_j(x_j) = \frac{h}{x_j} + \frac{h^2}{x_j^2}.$$

570 However, we noticed that for this last interval truncating the second term $\frac{h^2}{x_j^2}$, which is of order $O(h^2)$,
 571 improves the scheme's performance. We present both the performance of the original scheme and
 572 the truncated scheme in Table 2. Simulations for Table 2 are performed over the finite domain
 573 $x \in [0, 20]$.

574 *Example 4.* In this example, take $b(y, \cdot) = \frac{2}{y} dx$ and $a(x) = x^2$. Again, as given in [53], this
 575 problem has an exact solution of

576

$$\mu_t = (1 + 2t + 2tx) \exp(-x(1 + xt)) dx.$$

577 Numerical simulations are presented for this example in Figure 4 with $\Delta x = 1/40$. The BL error
 578 and order of convergence are presented in Table 3. Simulations for Table 3 are performed over the
 579 finite domain $x \in [0, 20]$.

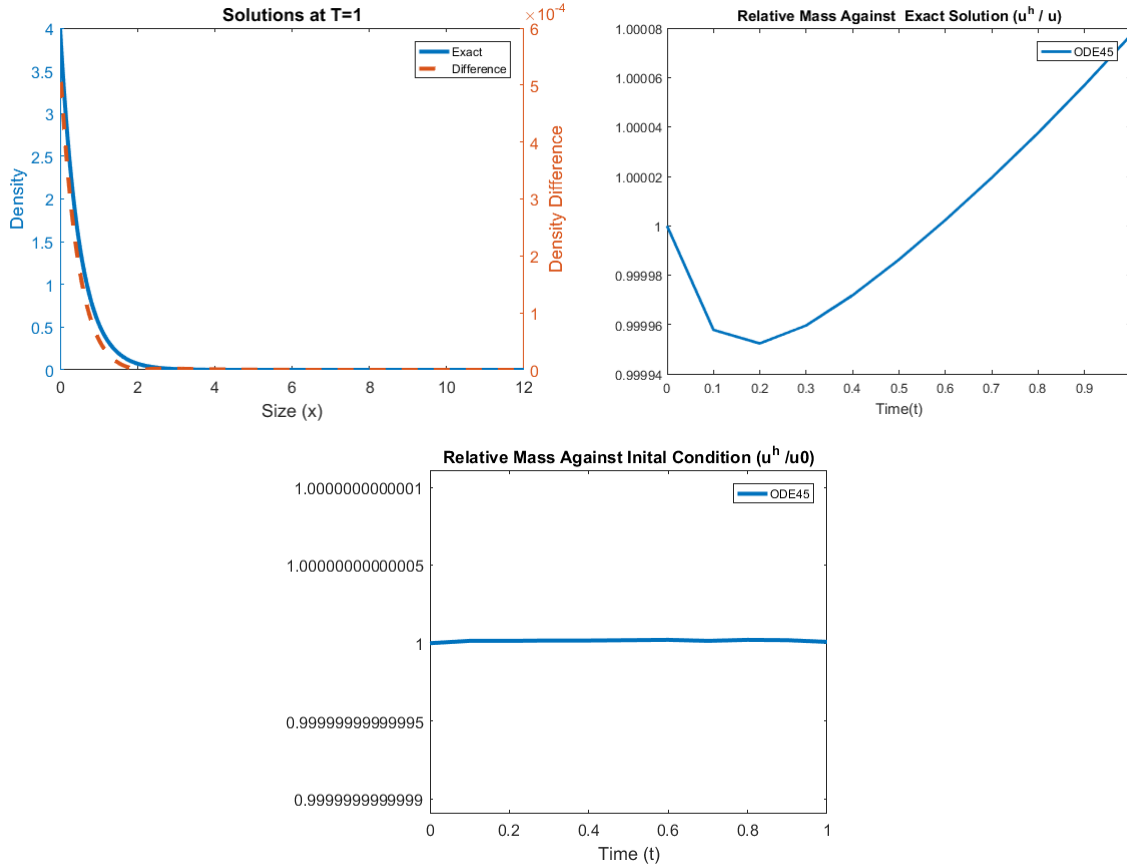


Figure 3: For Example 3 we present on the top left side the exact solution (solid line) and the absolute value of the difference between the exact and numerical solution (dashed-line). On the top right side we present the relative mass against the exact solution. On the bottom, we present the relative mass against the initial condition.

Number of Points	Original Scheme		Truncated Scheme	
	BL-Error	Order	BL-Error	Order
40	0.19243	NA	0.074275	NA
80	0.079672	1.2722	0.024212	1.6172
160	0.028642	1.4759	0.0068855	1.8141
320	0.0094434	1.6008	0.0018342	1.9084
640	0.0029433	1.6818	0.00047321	1.9546
1280	0.00088279	1.7373	0.00012017	1.9775

Table 2: Error and numerical order of convergence Example 3.

580 *Example 5.* For this example, we demonstrate the performance of the scheme for a domain where
581 the minimum size is positive. To this end, we truncate Example 3 above to the domain $[10^{-3}, 20]$
582 and use the kernel given by (6.4). Since the exact solution is not known for this equation, we
583 compare to the solution given in Example 3. Though we do not compute any numerical orders
584 of convergence, we point out the numerical and exact solutions in Figure 5 are very close. This

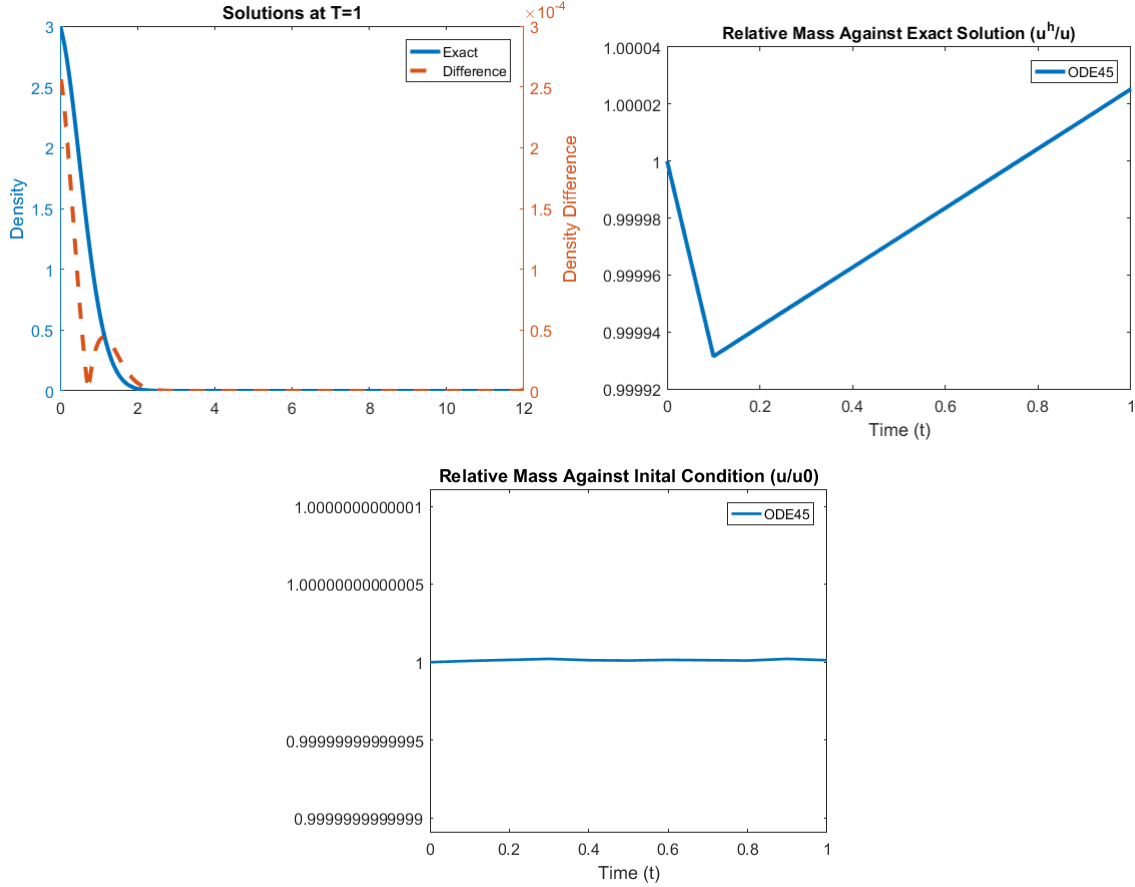


Figure 4: For Example 4 we present on the top left side the exact solution (solid line) and the absolute value of the difference between the exact and numerical solution (dashed-line). On the top right side we present the relative mass against the exact solution. On the bottom, we present the relative mass against the initial condition.

585 simulation is again done with $\Delta x = \frac{1}{40}$.

586 **Example 6.** In this example, we demonstrate what a discrete system would look like in our current
 587 frame work as well as provide an example of the results show in Theorem 5.1. We also demonstrate
 588 the mass conservation property of the coagulation terms of the scheme. The simulation is performed
 589 over the interval $[0, 20]$ however, for clarity we zoom into the interval $[0, 4]$. Take $k(x, y) \equiv 1$ and
 590 $\mu_0 = \delta_{0.2} + \delta_{0.4}$.

591 **7. Concluding Remarks.** In summary, we have presented a size-structured coagulation-fragmentation
 592 model formulated on the space of Radon measures endowed with the BL-norm. This model uni-
 593 fies the study of both the discrete and density based coagulation-fragmentation equations, both of
 594 which have been used in studying the dynamics of oceanic phytoplankton populations. We have
 595 shown, under biologically relevant assumptions, the model is well-posed using a fixed point approach
 596 discussed in recent papers [7, 8]. We also established a regularity result that shows, under certain
 597 conditions on the model parameters, the solution to the model is absolutely continuous to the left of
 598 the characteristic curve emanating from the point $(0, 0)$. This allows us to prove that any stationary
 599 solution of the model is absolutely continuous. This extends the result in [34] for structured popu-
 600 lation models without coagulation and fragmentation. Here, our proof differs from that in [34] since

Number of Points	Original Scheme		Truncated Scheme	
	BL-Error	Order	BL-Error	Order
40	0.1471		0.056501	NA
80	0.041762	1.8165	0.014505	1.9617
160	0.011112	1.9101	0.0036472	1.9917
320	0.0028655	1.9553	0.00091301	1.9981
640	0.00072752	1.9777	0.00022829	1.9998
1280	0.00018324	1.9893	0.000057021	2.0013

Table 3: Error and numerical order of convergence Example 4.

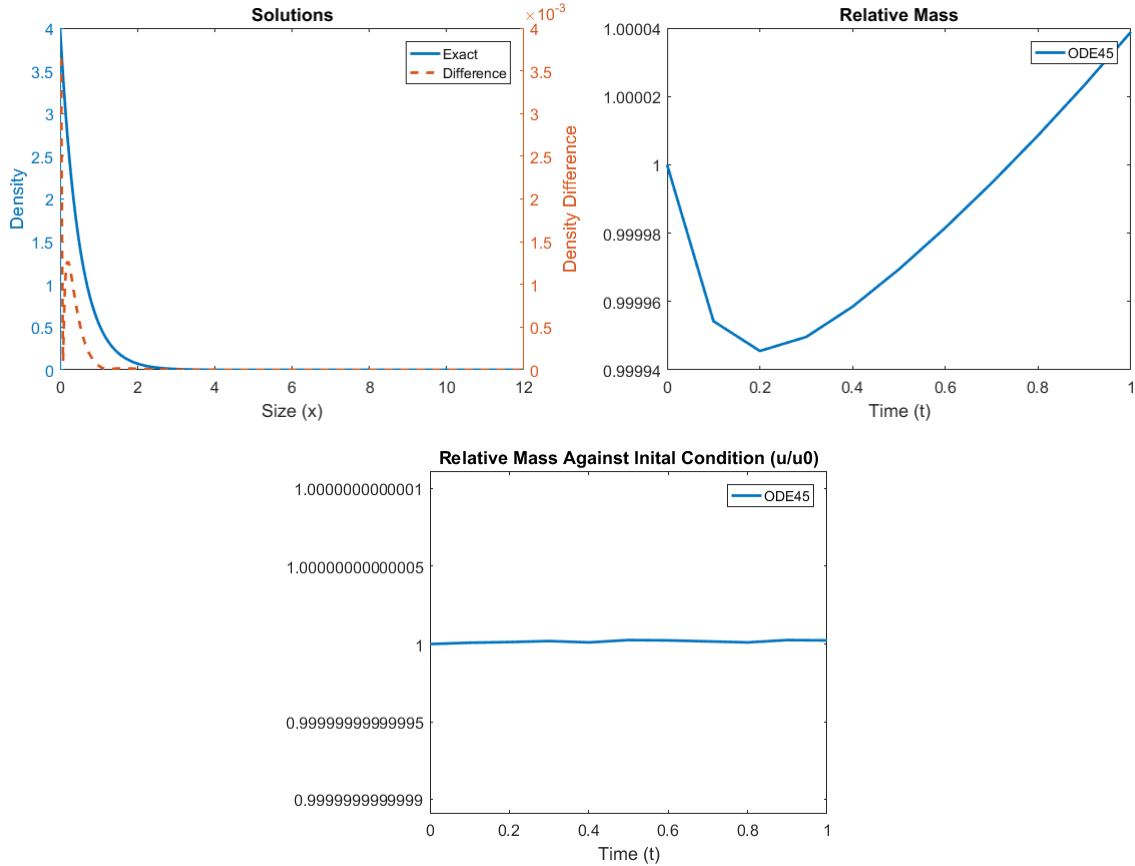


Figure 5: For Example 5 we present on the top left side the exact solution of Example 3 (solid line) and the absolute value of the difference between the exact and numerical solution (dashed-line). On the top right side we present the relative mass against the exact solution. On the bottom, we present the realtive mass against the initial condition.

601 it relies on the implicit fixed point representation of the measure valued solution. Furthermore,
602 we have shown how one obtains both the density and discrete coagulation-fragmentation equations
603 from model (3.1). We also provided a semidiscrete method for approximating solutions to these
604 equations and presented some numerical examples verifying our scheme. In these examples, we
605 observed the semidiscrete scheme appears to have at best a second order convergence rate in the BL

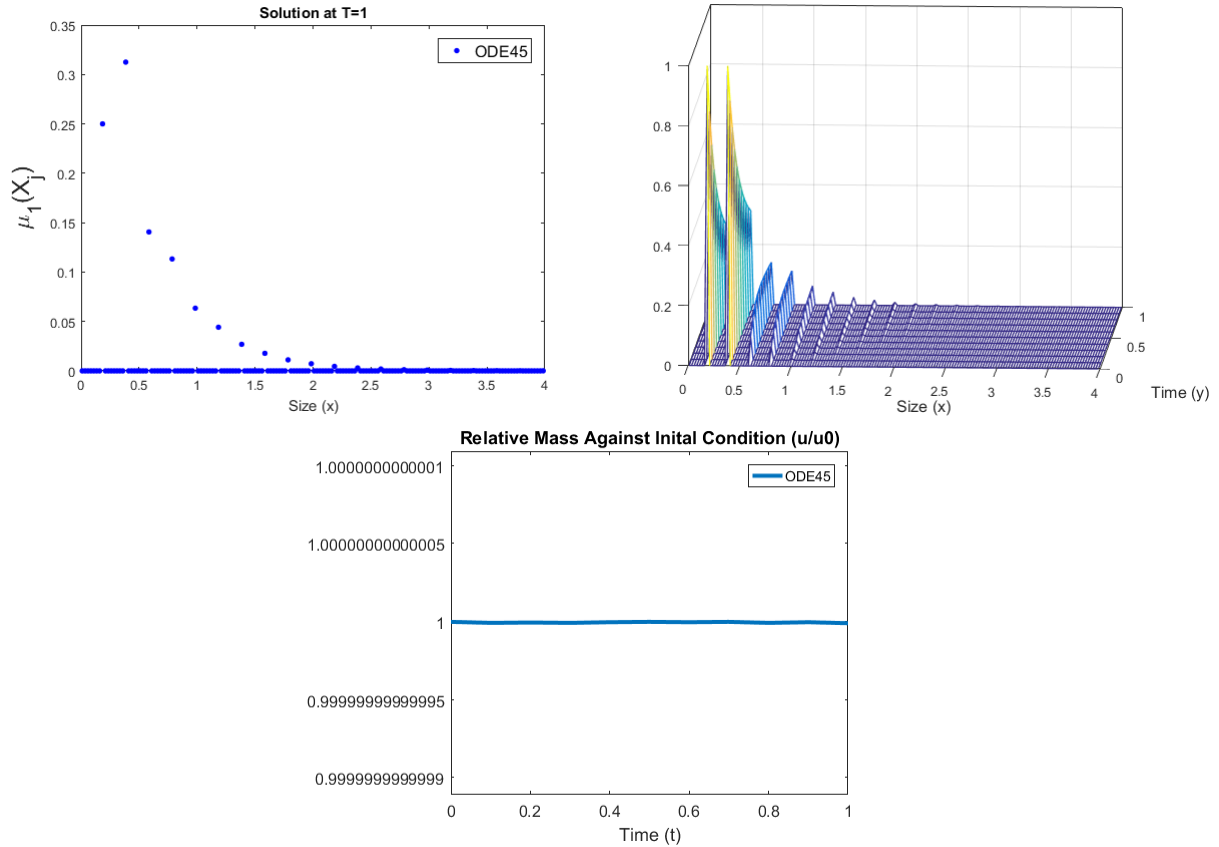


Figure 6: For Example 6 we present on the left side the numerical solution at time $T = 1$. On the right side we present a mesh of the solution over time. On the bottom, we present the relative mass according to the initial condition over $[0, 20]$.

606 norm. In addition to the cases covered by our convergence proof, the scheme also seems to perform
 607 well in the case of a gelation coagulation kernel.

608 While the semidiscrete scheme presented in this paper is convergent and conserves mass, it does
 609 not take into account a growth term. In the future, we plan to develop and study fully discrete
 610 higher order schemes for the full model (3.1) that preserves solution non-negativity and mass (e.g.
 611 [13, 45] in the space of integrable setting).

612

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