#### MIXED FINITE ELEMENTS FOR PLATES

# Ricardo G. Durán Universidad de Buenos Aires

- Necessity of 2D models.
- Reissner-Mindlin Equations.
- Finite Element Approximations.
- Locking.
- Mixed interpolation or reduced integration.
- General Error Analysis.
- Examples.

Why do we need to use 2D models if 3D elasticity can be solved by FE?

# 3D ELASTICITY EQUATIONS

 $D \subset \mathbb{R}^3$  Initial configuration of elastic solid.

$$u = (u_1, u_2, u_3)$$
 Displacement.

$$\varepsilon_{ij}(u) = \frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$
 Strain tensor.

$$2\mu \int_D \varepsilon(u) : \varepsilon(v) + \lambda \int_D \operatorname{div} u \operatorname{div} v = \int_D f \cdot v + \int_\Gamma g \cdot v$$
$$\forall v \in V \subset H^1(D)^3$$

Coercivity of this bilinear form follows from "Korn's inequality":

$$||u||_1 \le K||\varepsilon(u)||_0$$

Under appropriate boundary conditions.

**Remark 1.**: The constant K depends on the geometry of D.

#### RECALL CEA'S LEMMA:

If u is the exact solution and  $u_h$  the FE appromation then,

$$||u - u_h||_1 \le \frac{M}{\alpha} ||u - v_h||_1 \quad \forall v_h \in V_h$$

where M and  $\alpha$  are the continuity and coercivity constants of the bilinear form.

In our problem  $\alpha$  depends on the Korn's constant K. If K is too large then  $\alpha$  is too small and the constant in Cea's Lemma is large.

Consider a PLATE of thickness t:

$$D = \Omega \times (-t/2, t/2)$$

where  $\Omega \subset \mathbb{R}^2$  and t > 0, with t small in comparison with the dimensions of  $\Omega$ .

In this case:

$$K = O(t^{-1})$$

CONSEQUENCE:

THE METHOD IS NOT EFFICIENT: VERY SMALL MESH SIZE WILL BE NEEDED!

This is a serious computational drawback specially in 3D.

SOLUTION: USE 2D MODELS!

## REISSNER-MINDLIN EQUATIONS

$$\Omega \subset \mathbb{R}^2$$
  $D = \Omega \times (-t/2, t/2)$ 

Displacements are approximated by

$$u_1(x, y, z) \sim -z\theta_1(x, y)$$

$$u_2(x, y, z) \sim -z\theta_2(x, y)$$

$$u_3(x, y, z) \sim w(x, y)$$

 $\theta_1, \theta_2$  "rotations", w "transverse displacement".

Assuming a transverse load of the form

$$t^3 f(x,y)$$

and a that the plate is clamped,

$$\theta = (\theta_1, \theta_2)$$
 and  $w$ 

satisfy the system of equations:

$$a(\theta, \eta) + \kappa t^{-2} (\nabla w - \theta, \nabla v - \eta) = (f, v)$$
  
$$\forall \eta \in H_0^1(\Omega)^2, v \in H_0^1(\Omega).$$

$$a(\theta, \eta) := \frac{E}{12(1 - \nu^2)} \int_{\Omega} [(1 - \nu)\varepsilon(\theta)\varepsilon(\eta) + \nu \operatorname{div} \theta \operatorname{div} \eta],$$

E Young modulus,

 $\nu$  Poisson ratio,

 $\kappa := Ek/2(1+\nu)$  shear modulus,

k correction factor usually taken as 5/6.

We change notation and keep only the parameter t. So, our equations are

$$a(\theta, \eta) + t^{-2}(\nabla w - \theta, \nabla v - \eta) = (f, v)$$

For our purposes, the only important fact about a is that it is coercive in  $H_0^1$  (which follows from the 2D Korn inequality). We will not make other use of the explicit form o a.

The deformation energy is given by

$$\frac{1}{2}a(\theta,\theta) + \frac{t^{-2}}{2} \int_{\Omega} |\nabla w - \theta|^2 - \int_{\Omega} fw$$

It can be shown that the second term remains bounded when  $t \to 0$ . In particular,

$$t \to 0 \quad \Rightarrow \quad |\nabla w - \theta| \to 0$$

For the limit problem:

$$\nabla w = \theta$$
 "Kirchkoff constraint"

THIS IS A PROBLEM FOR THE NUMER-ICAL SOLUTION!

If t is small the problem is close to a constrained minimization problem.

FOR EXAMPLE: If we use standard  $\mathcal{P}_1$  finite elements for  $\theta$  and w, the restriction of the limit problem is too strong.

Indeed, if

$$\nabla w_h = \theta_h$$

then,  $\nabla w_h$  piecewise constant and continuous

$$\Rightarrow \nabla w_h$$
 constant

But,

$$\theta_h \in H_0^1(\Omega)^2 \quad \Rightarrow \nabla w_h = \theta_h = 0$$

CONSEQUENCE: For t small  $\theta_h, w_h \sim 0$ .

This is called "LOCKING"

Remark 2.: Indeed, now the continuity constant of the bilinear for is too large. It seems that we have a problem similar to the original 3D problem!

AND SO, WHAT IS THE ADVANTAGE OF USING THE 2D MODEL?

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SOLUTION: Mixed Interpolation or Reduced Integration.

IDEA: Relax the restriction of the limit problem.

 $\nabla w - \theta = 0$  replaced by  $\Pi(\nabla w - \theta) = 0$   $\Pi$  is some interpolation or projection onto some space  $\Gamma_h$ .

So, in the discrete problem, the restriction is verified only at some points or in some average sense.

FINITE ELEMENT APPROXIMATION:

$$\theta_h \in H_h \subset H_0^1(\Omega)^2, \quad w_h \in W_h \subset H_0^1(\Omega)$$
 are such that

$$a(\theta_h, \eta) + t^{-2}(\Pi(\nabla w_h - \theta_h), \Pi(\nabla v - \eta)) = (f, v)$$

$$\forall \eta \in H_h, v \in W_h$$

In the usual methods

$$\nabla W_h \subset \Gamma_h$$

So,

$$a(\theta_h, \eta) + t^{-2}(\nabla w_h - \Pi \theta_h, \nabla v - \Pi \eta) = (f, v)$$
$$\forall \eta \in H_h, v \in W_h$$

#### MIXED FORM:

Introducing the shear stress

$$\gamma = t^{-2}(\nabla w - \theta)$$

$$\begin{cases} a(\theta_h, \eta) + t^{-2}(\gamma_h, \nabla v - \Pi \eta) = (f, v) \\ \gamma_h = t^{-2}(\nabla w_h - \Pi \theta_h) \\ \forall \eta \in H_h, v \in W_h \end{cases}$$

MAIN PROBLEM: How to choose the spaces  $H_h$ ,  $W_h$ ,  $\Gamma_h$  and the operator  $\Pi$ ?

EXAMPLE: The Bathe-Dvorkin MITC4 rectangular elements (Mixed Interpolation Tensorial Components).

 $H_h$  and  $W_h$  are the standard  $\mathcal{Q}_1$  elements and  $\Gamma_h$  is locally defined by (a + by, c + dx)(is a rotated Raviart-Thomas space).

The operator  $\Pi$  is defined by

$$\int_{\ell} \Pi \eta \cdot t_{\ell} = \int_{\ell} \eta \cdot t_{\ell}$$

for all side  $\ell$  of an element, where  $t_{\ell}$  is the unit tangent vector on  $\ell$ .

# GENERAL ERROR ANALYSIS Continuous problem:

$$\begin{cases} a(\theta, \eta) + (\gamma, \nabla v - \Pi \eta) = (f, v) + (\gamma, \eta - \Pi \eta) \\ \gamma = t^{-2}(\nabla w - \theta) \\ \forall \eta \in H_0^1(\Omega)^2, v \in H_0^1(\Omega) \end{cases}$$

Discrete Problem:

$$\begin{cases} a(\theta_h, \eta) + (\gamma_h, \nabla v - \Pi \eta) = (f, v) \\ \gamma_h = t^{-2}(\nabla w_h - \Pi \theta_h) \\ \forall \eta \in H_h, v \in W_h \end{cases}$$

Error equation:

$$a(\theta - \theta_h, \eta) + (\gamma - \gamma_h, \nabla v - \Pi \eta)$$
  
=  $(\gamma, \eta - \Pi \eta) \quad \forall \eta \in H_h, v \in W_h$ 

Lemma 1. Let  $\theta_I \in H_h$ ,  $w_I \in W_h$  and  $\gamma_I = t^{-2}(\nabla w_I - \Pi \theta_I) \in \Gamma_h$ .

Suppose

$$\|\gamma - \Pi\gamma\|_0 \le Ch\|\gamma\|_1$$

and

$$(\gamma - \Pi \gamma, \eta) = 0 \quad \forall \eta \in \mathcal{P}_{k-2}^2$$

Let P be the  $L^2$  projection into  $\mathcal{P}_{k-2}^2$ . Then,

$$\begin{split} & \|\theta - \theta_h\|_1 + t\|\gamma - \gamma_h\|_0 \\ \leq & C(\|\theta_I - \theta\|_1 + t\|\gamma_I - \gamma\|_0 + h\|\gamma - P\gamma\|_0 \end{split}$$

Proof.

$$\begin{split} a(\theta_I - \theta_h, \eta) + (\gamma_I - \gamma_h, \nabla v - \Pi \eta) \\ = a(\theta_I - \theta, \eta) + (\gamma_I - \gamma, \nabla v - \Pi \eta) + (\gamma, \eta - \Pi \eta) \\ \forall \eta \in H_h, v \in W_h \end{split}$$

Take  $\eta = \theta_I - \theta_h$  and  $v = w_I - w_h$ . So,

$$\gamma_I - \gamma_h = t^{-2}(\nabla v - \Pi \eta)$$

Using the coercivity of a we obtain

$$\|\theta_I - \theta_h\|_1^2 + t^2 \|\gamma_I - \gamma_h\|_0^2$$

$$= a(\theta_I - \theta, \theta_I - \theta_h) + t^2(\gamma_I - \gamma, \gamma_I - \gamma_h) + (\gamma - P\gamma, \theta_I - \theta_h - \Pi(\theta_I - \theta_h))$$

and the lemma follows. •

To apply the Lemma we need to find approximations  $\theta_I$  and  $w_I$  such that the associated  $\gamma_I$  be also a good approximation. This will follow from the existence of approximations satisfying the following property

$$\boxed{\nabla w_I - \Pi \theta_I = \Pi(\nabla w - \theta)}$$

and

$$\|\theta - \theta_I\|_1 \le Ch^k \|\theta\|_{k+1}$$
$$\|\gamma - \Pi\gamma\|_0 \le Ch^k \|\gamma\|_k$$

This property can be seen as a generalization of the known Fortin property basic in the analysis of mixed methods. In fact, introducing the operators

$$I(\theta, w) = (\theta_I, w_I)$$
$$B(\theta, w) = \nabla w - \theta$$

and

$$B_h(\theta_h, w_h) = \nabla w_h - \Pi \theta$$

the property can be summarized by the following commutative diagram:

$$\begin{array}{ccc} (H_0^1)^2 \times H_0^1 & \xrightarrow{B} & (L^2)^2 \\ I \downarrow & & \downarrow \Pi \\ H_h \times W_h & \xrightarrow{B_h} & \Gamma_h \end{array}$$

When  $\Pi$  is an  $L^2$  projection, this is exactly the Fortin property, which is known to be equivalent to the inf-sup condition. This will be the case in one of our examples (the Arnold-Falk elements). In that case error estimates for the shear stress  $\gamma$  can also be obtained.

Remark 3. Indeed it is enough to have the commutative diagram with  $\Pi$  replaced by some operator with good approximation properties.

If approximations satisfying the properties above (i.e, commutative diagram + approximation properties) exist, then it follows from the Lemma:

#### ERROR ESTIMATES

$$\|\theta - \theta_h\|_1 + t\|\gamma - \gamma_h\|_0 + \|w - w_h\|_1$$

$$\leq Ch^k(\|\theta\|_{k+1} + t\|\gamma\|_k + \|\gamma\|_{k-1})$$

With constant C independent of h and t.

#### **EXAMPLES**

Example 1: MITC3 or DL element:

$$H_h: \mathcal{P}_1^2 \oplus \{\lambda_2\lambda_3t_1, \lambda_3\lambda_1t_2, \lambda_1\lambda_2t_3\}, \mathcal{C}^0$$

Degrees of freedom:  $\theta(V_i)$ ,  $\int_{\ell_i} \theta \cdot t_i$ 

$$W_h: \mathcal{P}_1 \mathcal{C}^0$$

$$\Gamma_h: (a-by,c+bx) \qquad \mathcal{C}^0 \text{ t. c.}$$

Degrees of freedom:  $\int_{\ell_i} \gamma \cdot t_i$ 

 $\Pi$  defined by

$$\int_{\ell} \Pi \gamma \cdot t = \int_{\ell} \gamma \cdot t$$

Example 2: Arnold-Falk element:

$$H_h: \mathcal{P}_1^2 \oplus \{\lambda_1\lambda_2\lambda_3\}, \mathcal{C}^0$$

Degrees of freedom:  $\theta(V_i)$ ,  $\int_T \theta$ 

$$W_h: \mathcal{P}_1$$
 Non Conforming

Degrees of freedom:  $\int_{\ell} w$ 

$$\Gamma_h: \qquad \mathcal{P}_0^2$$

 $\Pi$  defined by

$$\int_T \Pi \gamma = \int_T \gamma$$

 $L^2$ -projection!

Error estimates for Examples 1 and 2:

$$\|\theta - \theta_h\|_1 + \|w - w_h\|_1 \le Ch$$

Example 3: Bathe-Brezzi second order rectangular elements:

$$H_h: \mathcal{Q}_2, \mathcal{C}^0$$
 $W_h: \mathcal{Q}_2^r: \{1, x, y, xy, x^2, y^2, x^2y, xy^2\}, \mathcal{C}^0$ 
 $\Gamma_h: \{1, x, y, xy, y^2\} \times \{1, x, y, xy, x^2\}$ 
 $\Pi$  defined by

$$\begin{cases} \int_{\ell} \Pi \gamma \cdot t p_1 = \int_{\ell} \gamma \cdot t p_1 \\ \int_{R} \Pi \gamma = \int_{R} \gamma \end{cases}$$

Error estimates:

$$\|\theta - \theta_h\|_1 + \|w - w_h\|_1 \le Ch^2$$

Remark 4. Here the constant C is NOT independent of t because the right hand side involves higher order norms of the solution which depend on t.

### $L^2$ - ERROR ESTIMATES

$$\|\theta - \theta_h\|_0 + \|w - w_h\|_0 \le Ch^2$$

Recall the error equation

$$\begin{split} a(\theta-\theta_h,\eta) + (\gamma-\gamma_h,\nabla v - R\eta) &= (\gamma,\eta - R\eta) \\ \forall (\eta,v) \in H_h \times W_h, \end{split}$$

Duality argument:  $(\varphi, u) \in H_0^1(\Omega)^2 \times H_0^1(\Omega)$  solution of

$$\begin{cases} a(\eta,\varphi) + (\nabla v - \eta,\delta) \\ = (v,w-w_h) + (\eta,\theta-\theta_h) \\ \delta = t^{-2}(\nabla u - \varphi). \end{cases}$$
 
$$\forall (\eta,v) \in H_0^1(\Omega)^2 \times H_0^1(\Omega),$$

A priori estimate (we assume  $\Omega$  is convex)

$$\|\varphi\|_{2} + \|u\|_{2} + \|\delta\|_{0} + t \|\delta\|_{1}$$

$$\leq C(\|\theta - \theta_{h}\|_{0} + \|w - w_{h}\|_{0}),$$

Take  $v = w - w_h$  and  $\eta = \theta - \theta_h$  in the dual problem and use the error equation with  $(\eta, v) = (\varphi_I, u_I)$ :

$$\begin{split} \|w - w_h\|_0^2 + \|\theta - \theta_h\|_0^2 \\ = a(\theta - \theta_h, \varphi - \varphi_I) + t^2(\gamma - \gamma_h, \delta - \Pi\delta) \\ + (\theta_h - R\theta_h, \delta) + (\gamma, \varphi_I - \Pi\varphi_I), \end{split}$$

where we have used the commutative diagram property  $\Pi \delta = \frac{\kappa}{t^2} (\nabla \hat{u} - \Pi \hat{\varphi})$ .

PROBLEM: The last two terms.

For the MITC3 elements we have:

**Lemma 2.** If  $\gamma \in H(div, \Omega)$ ,  $\psi \in H_0^1(\Omega)$  and  $\psi_A$  is a piecewise-linear average interpolant:

$$|(\gamma, \psi_A - \Pi \psi_A)| \le Ch^2 ||div \gamma||_0 ||\psi||_1$$