

# Some optimization problems for nonlinear elastic membranes.

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$$\mathcal{J}(f) := \int_{\partial\Omega} f(x)u \, dS \longrightarrow \max$$

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- ▶  $u$  is the solution to

$$\begin{cases} -\Delta_p u + |u|^{p-2}u = 0 & \text{in } \Omega, \\ |\nabla u|^{p-2} \frac{\partial u}{\partial \nu} = f & \text{on } \partial\Omega. \end{cases}$$

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Observe that  $u$  is given by

$$\int_{\partial\Omega} fu \, dS = \sup_{v \in W^{1,p}(\Omega)} \frac{1}{p-1} \left\{ p \int_{\partial\Omega} fv \, dS - \int_{\Omega} |\nabla v|^p + |v|^p \, dx \right\}.$$

# History of the problem

- ▶ Related problem  $p=2$   $\rightarrow$  [Cherkaev - Cherkaeva, Adv. Math. Appl. Sci. 1999],
- ▶ Dirichlet BC and particular class  $\mathcal{A}$   $\rightarrow$  [Cuccu, Emamizadeh and Porru, EJDE 2006],

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Known results:

- ▶ Existence of optimal load  $f_*$ .
- ▶ Euler-Lagrange type equation for  $f_*$ .

# Applications

- ▶ Quasi-regular and quasi-conformal mappings in Riemannian manifolds with boundary
- ▶ Non-Newtonian fluids.
- ▶ Reaction–diffusion problems.
- ▶ Flow through porous media.
- ▶ Nonlinear elasticity.
- ▶ Glaciology.



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- ▶ Nonlinear elasticity.
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## Remark

Our results are new even in the linear case  $p = 2$ .

## Admissible classes

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- ▶ The second case  $\rightarrow$  EASY!
- ▶ The last case  $\rightarrow$  more interesting...

# The class of rearrangements

## Definition

We say that  $f$  is a rearrangement of  $f_0$  if

$$\mathcal{H}^{N-1}(f \geq t) = \mathcal{H}^{N-1}(f_0 \geq t) \quad t \geq 0.$$



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**Lemma** (Burton - McLeod, Proc. Roy. Soc. Edin. 1991)

Let  $f_0 \in L^q_+(\partial\Omega)$  and  $v \in L^{q'}_+(\partial\Omega)$ . Then there exists  $\hat{f} \in \mathcal{A}$  such that

$$\int_{\partial\Omega} \hat{f} v \, dS = \sup_{f \in \mathcal{A}} \int_{\partial\Omega} f v \, dS.$$

# The class of rearrangements

## Theorem

Let  $q > \frac{p'}{N}$ . There exists  $\hat{f} \in \mathcal{A}$  such that

$$\mathcal{J}(\hat{f}) = \int_{\partial\Omega} \hat{f} \hat{u} \, dS = \sup_{f \in \mathcal{A}} \mathcal{J}(f),$$

where  $\hat{u}$  is the solution to

$$\begin{cases} -\Delta_p \hat{u} + |\hat{u}|^{p-2} \hat{u} = 0 & \text{in } \Omega, \\ |\nabla \hat{u}|^{p-2} \frac{\partial \hat{u}}{\partial \nu} = \hat{f} & \text{on } \partial\Omega. \end{cases}$$

# The class of rearrangements

Proof.

Let

$$I = \sup_{f \in \mathcal{A}} \int_{\partial\Omega} f u_f dS,$$

where  $u_f$  is the solution associated to the load  $f$ .

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First, it is easy to see that  $I$  is finite. Then we take a maximizing sequence  $\{f_i\}_{i \geq 1}$  and let  $u_i = u_{f_i}$ .

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Now,  $u_i$  is bounded in  $W^{1,p}$ ,  $f_i$  is bounded in  $L^q(\partial\Omega)$  and so

$$u_i \rightharpoonup u \text{ weakly in } W^{1,p}(\Omega),$$

$$u_i \rightarrow u \text{ strongly in } L^p(\Omega),$$

$$u_i \rightarrow u \text{ strongly in } L^r(\partial\Omega), \quad r < p_* = \frac{(N-1)p}{N-p},$$

$$f_i \rightharpoonup f \text{ weakly in } L^q(\partial\Omega).$$

# The class of rearrangements

Proof.(cont.)

Then

$$\begin{aligned} I &= \lim_{i \rightarrow \infty} \int_{\partial\Omega} f_i u_i \, dS \\ &= \frac{1}{p-1} \lim_{i \rightarrow \infty} \left\{ p \int_{\partial\Omega} f_i u_i \, dS - \int_{\Omega} |\nabla u_i|^p + |u_i|^p \, dx \right\} \\ &\leq \frac{1}{p-1} \left\{ p \int_{\partial\Omega} f u \, dS - \int_{\Omega} |\nabla u|^p + |u|^p \, dx \right\}. \end{aligned}$$

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Now, by the Lemma of Burton – McLeod,

$$\int_{\partial\Omega} f u \, dS \leq \int_{\partial\Omega} \hat{f} u \, dS \quad \text{with } \hat{f} \in \mathcal{A}.$$



# The class of rearrangements

Proof.(the end)

Finally, denoting  $\hat{u} := u_{\hat{f}}$ , we get

$$\begin{aligned} I &\leq \frac{1}{p-1} \left\{ p \int_{\partial\Omega} f u \, dS - \int_{\Omega} |\nabla u|^p + |u|^p \, dx \right\} \\ &\leq \frac{1}{p-1} \left\{ p \int_{\partial\Omega} \hat{f} u \, dS - \int_{\Omega} |\nabla u|^p + |u|^p \, dx \right\} \\ &\leq \frac{1}{p-1} \left\{ p \int_{\partial\Omega} \hat{f} \hat{u} \, dS - \int_{\Omega} |\nabla \hat{u}|^p + |\hat{u}|^p \, dx \right\} \\ &= \int_{\partial\Omega} \hat{f} \hat{u} \, dS \\ &\leq I. \end{aligned}$$

This completes the proof.



## The unit ball of $L^q$

In this section, we consider the class

$$\mathcal{A} := \{f \in L^q(\partial\Omega) : \|f\|_{L^q(\partial\Omega)} = 1\}.$$

Again, in order for  $\mathcal{J}$  to make sense, we need  $q > \frac{p'}{N'}$ .

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$$S_{q'} = \min_{v \in W^{1,p}(\Omega)} \frac{\int_{\Omega} |\nabla v|^p + |v|^p dx}{\left( \int_{\partial\Omega} |v|^{q'} dS \right)^{p/q'}}$$

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Observe that the condition

$$q > \frac{p'}{N'} \leftrightarrow q' < p_* = \frac{(N-1)p}{N-p}$$

# The unit ball of $L^q$

More precisely,

$$\hat{f} = v_{q'}^{q'-1}, \quad \hat{u} = u_{\hat{f}} = \frac{1}{S_{q'}^{1/p-1}} v_{q'},$$

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where  $v_{q'}$  is a nonnegative, normalized extremal.  
In the linear setting,  $p = q = 2$ ,  $v_2$  is the first Steklov eigenfunction

$$\begin{cases} -\Delta v + v = 0 & \text{in } \Omega \\ \frac{\partial v}{\partial \nu} = \lambda v & \text{on } \partial\Omega. \end{cases}$$

# Characteristic functions

We consider the class

$$\mathcal{A} = \{\chi_D : D \subset \partial\Omega, \mathcal{H}^{N-1}(D) = A\}$$

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For technical reasons, it is better to work in the *relaxed* class

$$\mathcal{B} = \{\phi \in L^\infty(\partial\Omega) : 0 \leq \phi(x) \leq 1, \int_{\partial\Omega} \phi \, dS = A\}.$$



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It is well known that  $\mathcal{B}$  is the closure of  $\mathcal{A}$  with respect to the weak\* convergence in  $L^\infty$ .

# Characteristic functions

## Lemma

There exists  $\hat{\phi} \in \mathcal{B}$  such that

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Very similar to the rearrangement case (compactness argument). □

We like now to show that  $\hat{\phi}$  belongs to the class  $\mathcal{A}$ , i.e.

$$\hat{\phi} = \chi_{\hat{D}}, \quad \mathcal{H}^{N-1}(\hat{D}) = A.$$

# Characteristic functions

We need this result

## Lemma (Bathtub principle)

For any  $u \in L^p(\partial\Omega)$ , the maximization problem

$$I = \sup_{\phi \in \mathcal{B}} \int_{\partial\Omega} u(x)\phi(x) dS$$

is solved by

$$\hat{\phi}(x) = \chi_{\hat{D}}(x), \quad \{u > s\} \subset \hat{D} \subset \{u \geq s\}$$

where

$$s = \inf\{t : \mathcal{H}^{N-1}(\{u \geq t\}) \leq A\}.$$

The maximizer given is unique ( $\mathcal{H}^{N-1}$ -a.e.) if  $\mathcal{H}^{N-1}(\{u = s\}) = 0$ .

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This completes the proof of existence. □



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First step: **Shape derivative approach.**

## Shape Derivative

Assume that we fix a set  $D \subset \partial\Omega$  that is an open subset with regular boundary (i.e. its boundary relative to  $\partial\Omega$  is a smooth  $N - 2$  dimensional surface).

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We perturb  $D$  as follows:

- ▶ Take  $V: \mathbb{R}^N \rightarrow \mathbb{R}^N$  smooth, globally Lipschitz, with support in a neighborhood of  $\partial\Omega$  such that  $\langle V(x), \nu(x) \rangle = 0, \forall x \in \partial\Omega$ .

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- ▶ Define the flow

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- ▶ Define the perturbed domain  $D_t$  along the field  $V$  as

$$D_t = \psi_t(D)$$

(Observe that  $\psi_t(\Omega) = \Omega$  and  $\psi_t(\partial\Omega) = \partial\Omega$ )



# Shape Derivative

We denote by  $u_t$  the state of the system associated to the set  $D_t$ ,  
i.e.

$$\begin{cases} -\Delta_p u_t + |u_t|^{p-2} u_t = 0 & \text{in } \Omega \\ |\nabla u_t|^{p-2} \frac{\partial u_t}{\partial \nu} = \chi_{D_t} & \text{on } \partial\Omega \end{cases}$$

and define

$$j(t) = \mathcal{J}(\chi_{D_t}) = \int_{\partial\Omega} u_t \chi_{D_t} dS$$

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We want to compute

$$j'(0) = \frac{\partial \mathcal{J}}{\partial V}(D)$$

# Shape Derivative

## Theorem

$j(t)$  is differentiable at  $t = 0$  and we have

$$j'(0) = \frac{p}{p-1} \int_{\partial D} u_D \langle V, \nu_\tau \rangle d\mathcal{H}^{N-2}$$

where  $\nu_\tau$  stands for the exterior unit normal vector to  $D$  along  $\partial\Omega$ .

# Shape Derivative

Proof.

The proof follows from **lengthy** computations and adequate use of

- ▶ The *tangential* Changes of Variables Formula.

$$\int_{\Phi(\Gamma)} f \, dS = \int_{\Gamma} (f \circ \Phi) J_{\tau}(\Phi) \, dS$$

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- ▶ The *tangential* Divergence Theorem.

$$\int_D \operatorname{div}_{\tau} W \, dS = \int_{\partial D} \langle W, \nu_{\tau} \rangle \, d\mathcal{H}^{N-2} + \int_D H \langle W, \nu \rangle \, dS$$

( $\operatorname{div}_{\tau}(W)$  being the *tangential* divergence of  $W$ ,  $H$  is the mean curvature of  $D$ )

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Proof. (cont.)

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To overcome this problem, we work with the solutions to a regularized problem and then pass to the limit

$$\begin{cases} -\operatorname{div}(|\nabla u_t^\varepsilon|^2 + \varepsilon^2)^{(p-2)/2} \nabla u_t^\varepsilon + |u_t^\varepsilon|^{p-2} u_t^\varepsilon = 0 & \text{in } \Omega \\ (|\nabla u_t^\varepsilon|^2 + \varepsilon^2)^{(p-2)/2} \frac{\partial u_t^\varepsilon}{\partial \nu} = \chi_{D_t} & \text{on } \partial\Omega. \end{cases}$$



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$$\begin{cases} -\operatorname{div}((|\nabla u_t^\varepsilon|^2 + \varepsilon^2)^{(p-2)/2} \nabla u_t^\varepsilon) + |u_t^\varepsilon|^{p-2} u_t^\varepsilon = 0 & \text{in } \Omega \\ (|\nabla u_t^\varepsilon|^2 + \varepsilon^2)^{(p-2)/2} \frac{\partial u_t^\varepsilon}{\partial \nu} = \chi_{D_t} & \text{on } \partial\Omega. \end{cases}$$

Standard Elliptic Regularity Theory gives  $u_t^\varepsilon \in C^{2,\delta}$  and, moreover

$$u_t^\varepsilon \rightarrow u_t \quad \text{in } C^1 \quad (\varepsilon \rightarrow 0)$$



# Shape Derivative

By (much easier) similar computations we also get

Lemma

$$\frac{d}{dt} \mathcal{H}^{N-1}(D_t) |_{t=0} = \int_D \operatorname{div} V \, dS.$$

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Observe that, as  $V$  is tangential to  $\partial D$ , we can get

$$\int_D \operatorname{div} V \, dS = \int_D \operatorname{div}_\tau V \, dS = \int_{\partial D} \langle V, \nu_\tau \rangle \, d\mathcal{H}^{N-2}.$$

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Combining these formulas with the Lagrange Multipliers Rule, we finally obtain...

# Shape Derivative

## Corollary

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This is

$$\frac{p}{p-1} \int_{\partial D} u_D \langle V, \nu_T \rangle d\mathcal{H}^{N-2} = c \int_{\partial D} \langle V, \nu_T \rangle d\mathcal{H}^{N-2}$$

for every tangential vector field  $V$ . So the result follows. □



## Applications and open problems

- ▶ This type of *sufficient conditions* and *shape derivatives* have been used by many authors in developing numerical algorithms for computing optimal shapes. cf. E. Oudet, ESAIM COCV (2004) – Survey, JFB, P. Groisman and J.D. Rossi, Ann. Mat. Pura Appl. (2007), the book of O. Pironeau, Springer (1984), etc.

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THANK YOU